Astro 596/496 NPA Lecture 7 Sept. 9, 2009

Announcements:

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- Problem Set 1 due next time Note: intensity units in figure are 10^{-3} photons cm⁻² s⁻¹ sr⁻¹ for example: red $\rightarrow \sim 1.2 \times 10^{-3}$ photons cm⁻² s⁻¹ sr⁻¹
- TA Office Hours: 3-4pm Thursdays ... or by appointment

Last time: nuclear reactions and cross sections

Q: how is a cross section defined?

Q: what is the physical significance of σ ? units?

Q: how are cross sections related to particle sizes?

Cross Sections: Post-Holiday Reminder

For reaction $a + b \rightarrow$ something cross section *defined* as

$$\sigma_{ab} \equiv \frac{\text{reaction rate per target}}{\text{projectile flux}} = \frac{\Gamma_{\text{per }b}}{j_a} = \frac{\Gamma_{\text{per }a}}{j_b} \qquad (1)$$

where the flux of species *i* with speed *v* is
$$j_i = n_i v = \# \text{ particles area}^{-1} \text{ time}^{-1} \qquad (2)$$

Physically: cross section σ_{ab} is projected size of "bullseye" each reactant "sees" in the other \Rightarrow measures reaction probability/strength

Cross section controlled by physics of a + b interaction

- if short-ranged "contact" interaction: $\sigma \sim (\text{particle size})^2$
- \bullet if long-ranged interaction: σ unrelated to particle size

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 $Q: \sigma$ when distribution of speeds? most important example?

Reaction Rates for Thermal Particles

for $ab \to cd$, reaction rate per volume: $r_{ab\to cd} = n_a n_b \langle \sigma_{ab} v \rangle$ where $\langle \cdots \rangle$ is average over distribution of relative v

For non-relativistic gas at temperature T: relative speeds are in *Maxwell-Boltzmann* distribution

$$\langle \sigma v \rangle = \sqrt{\frac{8}{\pi \mu}} \frac{1}{(kT)^{3/2}} \int_0^\infty dE \ E \ \sigma(E) \ e^{-E/kT}$$
(3)

Note $e^{-E/kT}$: exponentially suppresses energies $E \gtrsim kT$

Recall Coulomb barrier:

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 $E_C = Z_1 Z_2 e^2/r = Z_1 Z_2$ 1.44 MeV (1 fm/r) Classically, need $E \gtrsim 1$ MeV to overcome barrier But kT = 0.86 keV $(T/10^7 \text{ K}) \ll E_C$ for solar-like temp !?!

Q: What does this seem to imply? Q: What's the flaw?

Quantum Mechanics to the Rescue

nuclei are quantum particles \rightarrow can tunnel Probability to tunnel under Coulomb barrier:

$$P \propto e^{-2\pi Z_1 Z_2 e^2/\hbar v} = e^{-bE^{-1/2}}$$

(4)

Also: geometrical factor: cross section $\sigma \propto \lambda_{deB}^2$, $\lambda_{deB} = \hbar/p$ de Broglie wavelength $\Rightarrow \sigma \propto 1/p^2 \propto 1/E$

expect σ functional form

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$$\sigma(E) = \frac{S(E)}{E} e^{-2\pi\eta} = \frac{S(E)}{E} e^{-bE^{-1/2}}$$

S(E): "astrophysical S-factor"

• S(E) encodes nuclear contribution to reaction

• S(E) often slowly varying with E

Q: if so, σ behavior at large E? small E?

• σ and S-factor for ${}^{3}\text{He}(\alpha,\gamma){}^{7}\text{Be}$

Thermonuclear Rates

So: thermonuclear rates reduce to:

$$\langle \sigma v \rangle = \langle \sigma v \rangle_T$$

= $\sqrt{\frac{8}{\pi \mu}} \frac{1}{(kT)^{3/2}} \int_0^\infty dE \ S(E) e^{-E/kT - bE^{-1/2}}$

Procedure:

(1) astro theory/obs identifies needed reaction

- (2) nuclear expt: measure $\sigma(E) \to S(E)$
- (3) find $\langle \sigma v \rangle$ vs T (usually numerically)
- (4) fit result to function

 $^{\circ}$ Q: note exponential-behavior vs E? implications?

The Gamow Peak

integrand $S(E)e^{-G(E)}$ peaks at/near *minimum* of exponential $G(E) = E/kT + bE^{-1/2}$ min at G' = 0: "most effective energy" or "*Gamow Peak*" $E_0 = (bkT/2)^{2/3}$, where

$$G_{\min} \equiv \tau = G(E_0) = 3(b^2/4kT)^{1/3}$$
(5)
= $4.25(Z_1^2 Z_2^2 A)^{2/3} \left(\frac{10^9 \text{ K}}{T}\right)^{1/3}$ (6)

Q: behavior with *T*? interpretation?

expand exponential around peak energy E_0 :

$$G(E) \approx \tau + \frac{1}{2} G''(E_0) \left(E - E_0 \right)^2$$
 (7)

σ

use expansion of exponential $G(E) \approx \tau + \frac{1}{2}G''(E_0) (E - E_0)^2$ in thermonuclear integral (method of steepest descent)

Then we have

$$\begin{array}{lll} \langle \sigma v \rangle &\simeq & \sqrt{\frac{8}{\pi \mu}} \frac{1}{(kT)^{3/2}} S(E_0) e^{-\tau} \int_{-\infty}^{\infty} dE \ e^{-(E-E_0)^2/2\Delta^2} \\ &= & \frac{8}{9\sqrt{3}\pi Z_1 Z_2 e^2 m} \tau^2 e^{-\tau} S(E_0) \\ &\propto & T^{-2/3} e^{-a/T^{1/3}} \end{array}$$

Q: behavior at high T? low T? are these reasonable?

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Mean Lifetimes

for reaction $i + j \rightarrow k + l$ define mean lifetime $\tau_i(ij)$ of i against reaction with j as

$$(\dot{n}_i)_{ij} = -\frac{n_i}{\tau_i(ij)} \tag{8}$$

or $\tau_i(ij) = \|n_i/\dot{n}_i\|$

But
$$\dot{n}_i = -r_{ij} = -n_i n_j \langle \sigma v \rangle_{ij}$$

 $\Rightarrow \tau_i(ij) = 1/n_j \langle \sigma v \rangle_{ij} = 1/\Gamma_{\text{per}\,i}(ij)$

useful to write

$$\Gamma_{\text{per}\,i}(ij) = n_j \langle \sigma v \rangle_{ij} \simeq \frac{X_j}{A_j} \frac{\rho}{m_u} \langle \sigma v \rangle_{ij} = \frac{X_j}{A_j} \rho[ij] \tag{9}$$

where $[ij] = \langle \sigma v \rangle_{ij} / m_u = N_{Avo} \langle \sigma v \rangle_{ij}$ given in tabulations

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Why would this be a useful form?

Partial Lifetimes: Examples

Reactions in the Sun

In the solar core today: $T \simeq 16 \text{ MK} = 1.6 \times 10^7 \text{ K}$ density $\rho \simeq 150 \text{ g cm}^{-3}$ $X_{\text{H}} \simeq 0.33$

What is lifetime of a deuteron against $d(p,\gamma)^3$ He?

$$\frac{1}{\tau_d(pd)} = X_{\rm H}\rho[dp \to \gamma^3 {\rm He}]$$

$$\simeq (0.3)(150 \,{\rm g \, cm^{-3}})(2 \times 10^{-10} \,{\rm cm^3 \, s^{-1} \, g^{-1}})$$

$$\sim 10^{-8} \,{\rm s^{-1}}$$

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or $\tau_p(pd) \sim 3$ yrs: "immediately"

compare ${}^{16}O(p,\gamma){}^{17}F$ Exponential factor $\tau(E_0) \sim 4 \times$ larger! $\tau_O(p{}^{16}O) \sim 10{}^{57} \text{ s} \gg \text{age of Univ!}$ no ${}^{16}O$ burned in solar core (on main seq)

Particle Physics

Antimatter

Fundamental result of Relativistic QM:

every particle has an antiparticle

- $\overline{e^-} = e^+$ positron
- $\bar{p} = antiproton$

Fermilab: $p\bar{p}$ collisions

antimatter is **not** second class citizen! e.g.: e⁺ totally *stable* when left alone *So why so volatile in the lab?*

www: e^+ annihilation in Galactic center

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