Astro 596/496 NPA Lecture 39 Dec. 2, 2009

Announcements:

• Problem Set 6 due

Last time: enlarged view of "multi-messenger" astrophysics i.e., cosmic probes beyond photons

- neutrinos
- dark matter (?)
- gravity waves
- cosmic rays
- Cosmic Rays: Q: what are they?
 Q: how do they propagate?

Propagation and Cosmic-Ray Abundances

Consider "primary" cosmic ray species abundant at source, e.g., Fe

eventually lost from CR, in one of two ways:

• escape

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• fragmentation = "spallation" in interaction with interstellar gas, e.g., $Fe_{cr} + p_{ism} \rightarrow Mn + \cdots$

in practice, escape dominates loss but spallation not negligible

Q: how will spallation affect CR abundances?

Spallation and Cosmic-Ray Propagation

spallation "erodes" all primary species especially the most abundant

• for primary species, spallation is a (small) *sink*

but also produces fragments, typically a few nucleons lighter

• for these "secondary" nuclei, spallation is source!

Net effect of spallation

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- reduce CR abundance "peaks"
- and "fill in" CR abundance "valleys"

origin of CR abundance differences with solar!

www: cosmic-ray vs solar abundances

The Leaky Box Model

Realistic CR propagation controlled by magnetic ISM $\vec{B}(\vec{x})$ inhomogeneities \rightarrow CR *random walk*, jump field lines \rightarrow cosmic-ray propagation is **diffusive**

full propagation equation complex: diffusion, advection, (re)acceleration energy losses due to ionization, collisions

but to see main effects on spectrum, can simplify: diffusion term \rightarrow timescale for escape from Galaxy "leaky box" model

Assume isotropic: net particle flux $\phi(\varepsilon) = \int d\Omega I(\varepsilon) = 4\pi I(\varepsilon)$

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For
$$n_{\varepsilon} = dn/d\varepsilon = \phi(\varepsilon)/v(\varepsilon)$$
, leaky box:

$$\frac{\partial n_{\varepsilon}}{\partial t} = -\text{energy loss} - \text{collisions} - \text{escape} + \text{sources}$$

$$= -\frac{\partial}{\partial \varepsilon} (bn_{\varepsilon}) - n_{\text{ISM}} \sigma_{\text{inel}} vn_{\varepsilon} - \frac{n_{\varepsilon}}{\tau_{\text{esc}}} + q_{\varepsilon}$$

where $b = |\partial E / \partial t|$ is rate of energy loss \rightarrow ISM ionization

define matter "thickness" $dX = \rho_{\text{ISM}} dx = \rho_{\text{ISM}} v dt$

$$\frac{\partial \phi}{\partial X} = -\frac{\partial}{\partial \varepsilon} (w\phi) - \frac{\phi}{\Lambda_{\text{inel}}} - \frac{\phi}{\Lambda_{\text{esc}}} + Q_{\varepsilon}$$
(1)
where $w = b/\rho v$, $\Lambda_{\text{inel}} = m/\sigma_{\text{inel}}$, $\Lambda_{\text{esc}} = \rho v \tau_{\text{esc}}$, $Q = q/\rho$
 $[\Lambda] = g/\text{cm}^2$ "grammage"

for high ε , ioniz. scale $\varepsilon/w \ll \Lambda$: neglect then in steady state: $\partial/\partial t = \partial/\partial X = 0$

 $\phi = \Lambda Q = q\tau$: flux = sources × escape time

CR Sources: primary (present at source): $Q_i \simeq Q_{accel}$ secondary (made in flight): $Q_\ell \simeq Q_{spall} = n_j \sigma_{ij}^\ell \phi_i / \rho$

From secondary/primary ratio
$$\phi_{\ell}/\phi_i \simeq \Lambda n_j \sigma_{ij}^{\ell}/\rho$$

find $\Lambda_{\rm esc} \simeq 10 \text{ g/cm}^2 \ll \Lambda_{\rm inel}$
 $\Rightarrow \tau_{\rm esc} \sim \Lambda/m_p \langle n \rangle c \simeq 6 \times 10^6 \text{ yr}$
using $\langle n \rangle = \langle n \rangle_{\rm ISM} = 1 \text{ cm}^{-3}$

Note: can make radioactive secondaries e.g., ¹⁰Be ($\tau = 1.5$ Myr) ep "clock" \rightarrow CR age $t \sim 2 \times 10^7$ yr higher than naive τ_{esc} interpretation: $\langle n \rangle \sim 1/3 \langle n \rangle_{ISM}$ \Rightarrow CR spend part of lives in low-density Galactic halo

Acceleration

How does nature accelerate particles?

Hint: in solar system, see low-E particle accel.:

- in coronal mass ejections
- in planetary bowshocks
- at the solar wind termination shock
- ⇒ at magnetized, collisionless shocks

Magnetized

- charged particles spiral around field lines speed $v_{||}$ along (parallel to) field \ll total speed v
- \Rightarrow charged particles carried ("entrained")

in flow between scatterings

Shocks

- discontinuities in gas flows
- deceleration: supersonic \rightarrow subsonic
- ordered pre-shock ("upstream") flow converted to disordered, turbulent post-shock downstream flow
- \Rightarrow downstream magnetic field tangled \rightarrow scattering centers

Collisionless

charged particles scatter off magnetic inhomogeneities, but don't collide with (and lose energy to) gas particles

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Q: how would such shocks lead to acceleration?

Diffusive Shock Acceleration

basic idea: 1st-order Fermi mechanism:

- incoming (upstream) particles encounter shock
- scatter collisionlessly: bounce between slower downstream and faster upstream flows with some probability of escape (advection) downstream
- like ball between convering walls: energy gain each cycle

net effect: power law spectrum!

www: simulation of diffusively accelerated particle trajectory

In general: astrophysical shocks are collisionless and magnetized, so astrophysical shocks accelerate particles stronger shocks \rightarrow more powerful/higher-E acceleration

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Q: candidates for cosmic-ray acceleration in Milky Way?

Cosmic Ray Accelerator Candidates

Leading candidate for CR acceleration: shock acceleration in supernova blast waves

- strong shocks
- large energy reservoir
- long-lasting

www: e (and ion?) acceleration to TeV in SN 1006

The Energetics & Origin of Cosmic Rays

Local: energy densities

 $\epsilon_{CR} \sim \epsilon_{therm,ISM} (= n_{ISM} kT) \sim \epsilon_{mag} (= B^2/8\pi) \sim \epsilon_{CMB} \sim 1 \text{ eV cm}^{-3}$ Q: why is CMB energy density comparable to the rest? Q: why is mag energy comparable to CR?

Global: cosmic ray sources and escape if steady state, $L_{accel} = L_{esc}$, but

$$L_{esc} = \frac{N_{CR}}{\tau_{esc}}$$
$$= \frac{\frac{E_{CR}n_{CR}V_{CR}}}{\Lambda/\rho_{ism}v}$$
$$= \frac{\frac{E_{CR}\Phi_{CR}M_{gas}}}{\Lambda} \sim 10^{41} \text{ erg/s} \sim 0.3 \text{ foe/century}$$

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If SN rate is \sim 3/century, then need \sim 10% of E into particles