

Astro 596/496 PC

Lecture 36

April 21, 2010

Announcements:

- PS6 out, due in class *Wednesday* next week
- **ICES** available online – please do it!

Last time: began CMB anisotropies

Q: dark matter perturbations before matter-rad eq? after?

Q: baryonic perturbations before recomb? after?

Q: physical origin of CMB ΔT on scales $\gg d_{\text{hor, recomb}}$?

dark matter: pressureless

→ all k modes unstable if inside Hubble length

but: perturbations grow verry sloooowly during radiation era

→ dark matter structures begin formation at matter-radiation equality

baryons: until recomb, tightly coupled to photons

→ feel huge photon pressure $P_\gamma \propto T^4$

→ sound speed $c_s \sim c/\sqrt{3}$ huge!

so all sub-horizon modes stable! just oscillate

→ relativistic pressure-mediated (i.e., acoustic) standing waves!

oscillation frequency $\nu = c_s/\lambda$:

small-scale modes oscillate many times

↳ largest-scale modes $\lambda = c_s\eta_{\text{hor}}$ oscillates only once

Inflation and Sachs-Wolfe

Inflation: quantum fluctuations \rightarrow density fluctuations

- adiabatic (all species)
- Gaussian
- scale invariant—what does this mean?

In detail: inflation predicts that the dimensionless fluctuations in the *gravitational potential* \leftrightarrow *local curvature* are independent of scale

\rightarrow this was what we really calculated in Inflation discussion

inflationary scale-invariance is for grav potential:

i.e., Fourier mode contribution $\Delta_{\Phi}^2 \sim k^3 |\Phi_k|^2 \sim \text{const}$ indep of k

ω
 \rightarrow scale invariant: $|\Phi_k|^2 \sim k^{-3}$

need to connect gravitational potential/curvature perturbations to density perturbations

But in Newtonian regime, know how to do this:

Poisson relates potential and density:

$$\nabla^2 \delta\Phi = 4\pi G \delta\rho \rightarrow \Phi_k \sim \delta_k / k^2$$

$$\text{and so } P(k) = |\delta_k|^2 \sim k^4 |\Phi_k|^2$$

thus scale invariant gravitational potential gives power spectrum:

$$P(k) \sim k^4 |\Phi_k|^2 \sim k \tag{1}$$

i.e., scale invariance: $P(k) \sim k^n$, $n_{\text{scale-inv}} = 1$

Angular vs Linear Scales

So far: decomposed fluctuations in (3-D) \vec{k} -space
but observe on sky: 2-D angular distribution

Transformation: projection of plane waves
at fixed k : see intersection of wave with last scattering shell
www: Wayne Hu animation

appears on a range of angular scales
but typical angular size is $\theta \sim \lambda/d_{\text{rec,com}} \sim (kd_{\text{rec,com}})^{-1}$

large angles \rightarrow large λ (check!)

for large angular scales $\theta > \theta_{\text{hor,diam}} \sim 1^\circ$, superhorizon
perturbations not affected by oscillation

for small angular scales, see standing waves

- peaks at extrema, harmonics of sound horizon
 k are in ratios 1:2:3:...

The CMB Observed

- Observe 2-D sky distribution of $\frac{\Delta T}{T}(\hat{n}) \equiv \Theta(\hat{n})$ in direction \hat{n}
- Decompose into spherical harmonics

$$\Theta(\hat{n}) = \sum_{\ell=2}^{\infty} \sum_{m=-\ell}^{\ell} a_{\ell m} Y_{\ell m}(\theta, \phi) \quad (2)$$

with $Y_{\ell m}$ spherical harmonics Q: why not $\ell = 0, 1$?

Q: angular size vs ℓ ? λ vs ℓ ?

Form angular correlation function Q: what is this physically?

$$\langle \Theta(\hat{n}_1) \Theta(\hat{n}_2) \rangle = \frac{1}{4\pi} \sum_{\ell=2}^{\infty} (2\ell + 1) \langle |a_{\ell m}|^2 \rangle P_{\ell}(\hat{n}_1 \cdot \hat{n}_2) \quad (3)$$

$$= \frac{1}{4\pi} \sum_{\ell=2}^{\infty} (2\ell + 1) C_{\ell} P_{\ell}(\cos \vartheta) \quad (4)$$

° where $\cos \vartheta = \hat{n}_1 \cdot \hat{n}_2$

Q: averaged over the m azimuthal modes—why?

all interesting anisotropy information encoded in

$$C_\ell = \langle |a_{\ell m}|^2 \rangle \quad (5)$$

isotropy \rightarrow azimuthal dependence averages to zero

Note: analog of Δ^2 (variance per log scale) is
 $\mathcal{T}^2(\ell) = \ell(\ell + 1)C_\ell$: usually what is plotted

Since $P_\ell(\cos \theta) \sim (\cos \theta)^\ell \sim \cos(\ell\theta)$
at fixed ℓ , angular size $\theta \sim 2\pi/\ell = 180^\circ/\ell$
e.g., $\ell = 2$ quadrupole $\rightarrow \theta \sim 90^\circ$
and horizon size $\theta \sim 1^\circ$ is at $\ell \sim 200$

and since $\theta \sim \lambda/d_{\text{rec}} \sim 1/dk$:

\sim multipoles scale as $\ell \sim 1/\theta \sim k \sim 1/\lambda$

low $\ell \rightarrow$ big angular, physical scales \rightarrow small k

CMB Anisotropy Observations: Strategy

- achieve high sensitivity, remove systematics
make a “difference experiment”
i.e., measure δT directly, don't subtract
- observe as much of the sky as possible (or as needed!)
balloons/ground: limited coverage
satellites (COBE, WMAP, Planck): all-sky
- remove Galactic contamination: “mask” plane
- recover Θ for observed region

- decompose into spherical harmonics $Y_{\ell m}$
- construct power spectrum $\ell(\ell + 1)C_\ell$
- report results
- collect thousands of citations, prominent Prizes

CMB Temperature Anisotropies: Results

COBE (1993)

- first detection of $\delta T/T \neq 0$
- receiver horn angular opening $\sim 8^\circ$
→ only sensitive to large angular scales
i.e., superhorizon size
- found $(\delta T/T)_{\text{rms}} \sim 10^{-5}$
- power $\ell(\ell + 1)C_\ell$ flat → implies $P(k) \sim k!$
 $n = 1$ spectrum: scale invariant!

Interregnum (late 90's, early 00's)

- ground-based, balloons confirmed anisotropy
- acoustic peaks discovered
strong indication of first peak

WMAP (2003-)

- first all-sky survey of small angular scales
- $n = 1$ confirmed, indication of small tilt $n - 1 \neq 0$?
consistent with inflation! and non-trivially so!
- acoustic peaks mapped: good measurement of 1st, 2nd
detection of third
- first peak: $\ell \sim 200$ horizon at recomb!
- power dropoff seen at large ℓ
→ nonzero thickness of last scattering
due to photon diffusion, non-instantaneous decoupling