

Astro 210  
Lecture 9  
Feb 7, 2011

Announcements

- HW3 due Friday  
HW1 Q8 bonus still available
- iClicker scores to date are posted
- Planetarium: new shows added  
registration, report form, info online

# Homework, Collaboration, and Academic Honesty

From the course Syllabus, emphasis added:

Discussing course material with your classmates is in general not only allowed but in fact a good idea. However, each student is expected to do *his or her own work*. On homework, you may discuss the questions and issues behind them, but *you are responsible for your own answers*.

It should be obvious but:

you are *not* to consult homework solutions from previous semesters

## Last Time: Sir Isaac Weighs In

Newton's Laws of motion

I. inertia *Q: just a special case of Newton II?*

II.  $\vec{F} = m\vec{a}$  *Q: fortunetelling & archæology?*

III. action-reaction

*Q: when/where/to what do Newton's laws of motion apply?*

Newtonian Gravitation

*Q: magnitude of gravity force between masses  $m$  and  $M$   
with distance  $r$ ?*

*Q: direction of the force?*

# Kepler from Newton

## Kepler I: Orbits are ellipses

Newton: bound orbits due to gravity are ellipses: check!

## Kepler II: Equal areas in equal times

Newton: consider small time interval  $dt$

move angle  $d\theta = \omega dt$

sweep area

diagram: top view: path,  $d\theta$ ,  $\vec{r}$ ,  $\vec{v}$ ,  $\vec{v}_\theta$

$$dA = \frac{1}{2}r^2 d\theta = \frac{1}{2}r^2 \omega dt \quad (1)$$

but  $\omega = v_\theta/r$ , where  $\vec{v}_\theta \perp \vec{r}$

$\Rightarrow$  swept area

+

$$dA = \frac{1}{2}r^2 \frac{v_\theta}{r} dt = \frac{1}{2}r v_\theta dt \quad (2)$$

⇒ swept area

$$dA = \frac{1}{2} r^2 \frac{v_\theta}{r} dt = \frac{1}{2} r v_\theta dt \quad (3)$$

finally,  $r v_\theta = |\vec{r} \times \vec{v}| = |\vec{L}|/m$

Q: *why?*, so

$$dA = \frac{1}{2} \frac{L}{m} dt \quad (4)$$

Woo hoo! were' home free! Q: *why?*

But  $L = \text{const}$  for radial force ( $\vec{r} \times \vec{F} = 0$ )

so

$$\frac{dA}{dt} = \frac{L}{2m} = \text{const} \quad (5)$$

Kepler II!  $\rightarrow$  comes from ang. mom. cons.!

## iClicker Poll: Orbits and Angular Momentum

Consider two objects at 1 AU:

- a planet in a circular orbit
- a spacecraft moving directly towards the Sun

Which has the larger *magnitude* of angular momentum?

- A** the planet
- B** the spacecraft
- C** either, depending on the speed of the spacecraft

**Kepler III:**  $a^3 = kP^2$

Newton: can prove generally for elliptical orbits  
bad news: price is lotsa algebra

good news: simple to do for circular orbits  
circular  $\rightarrow r = a$ , and  $v^2 = GM/a$   
but also  $v = 2\pi a/P$  Q: why?

$$v^2 = \left(\frac{2\pi a}{P}\right)^2 = \frac{4\pi^2 a^2}{P^2} \quad (6)$$

$$= \frac{GM}{a} \quad (7)$$

$$\Rightarrow a^3 = \left(\frac{GM}{4\pi^2}\right) P^2 \quad (8)$$

check!

$\infty$  bonus:  $k = GM/4\pi^2$  depends on mass of central object  
 $\rightarrow$  same  $k$  for all planets

# Energy

For “test” particle  $m$  moving due to gravity of  $M$   
Gravitational potential energy: Q: why “potential”?

$$PE = -GMm/r$$

Kinetic energy:

$$KE = \frac{1}{2}m\vec{v}^2 = \frac{1}{2}m(v_x^2 + v_y^2 + v_z^2) = \frac{1}{2}m\dot{r}^2 \quad (9)$$

Total energy:

$$TE = KE + PE = -GMm/r + \frac{1}{2}m\dot{r}^2$$

key result:  $d(TE)/dt = 0$

→ total energy conserved!

◦ that is: value of  $TE$  the same for all time!

## Orbits Revisited

Bound orbits (circle & ellipse): in polar coordinates

$$r(\theta) = \frac{(1 - e^2)a}{1 + e \cos \theta} \quad (10)$$

**Circle** radius  $r = a = \text{const}$ , eccentricity  $e = 0$

recall: circular orbit has constant speed  $v_C^2 = GM/r$

$$PE = -\frac{GMm}{r} < 0 \quad (11)$$

$$KE = \frac{1}{2}mv_C^2 = \frac{1}{2}m\frac{GM}{r} = \frac{1}{2}\frac{GMm}{r} = -\frac{1}{2}PE \quad (12)$$

$$\Rightarrow TE = KE + PE = PE/2 = -|PE|/2 < 0 \quad (13)$$

10  $TE < 0$ : negative? yes!

Q: what does it mean to have negative energy?

for orbiting system  $TE < 0$ :

→ have to *supply* energy to system to break it apart

Why? when particles are *at rest* and “very” *far apart*

$$KE = mv^2/2 = 0$$

$$PE = GMm/r \rightarrow 0 \quad Q: \text{how far apart is this?}$$

and so  $TE = KE + PE = 0$ : zero total energy

But if start in closed orbits (circular or elliptical):  $TE < 0$

→ To “break” the system from closed orbits, must *add* energy

But energy is conserved → not spontaneously added

so system is **bound**

⇒ can't fall apart without external influence

Note:  $KE = -PE/2 = |PE|/2$  generally true for

≡ gravitating systems in equilibrium:

**“virial theorem”**

**ellipse**: semimajor axis  $a$ , eccentricity  $0 < e < 1$

turns out:  $TE$  depends only on  $a$ , not  $e$

from cons of energy

$$TE = -GMm/r + \frac{1}{2}mv^2 = -GMm/2a < 0 \rightarrow \text{bound}$$

can show

$$v^2 = GM \left( \frac{2}{r} - \frac{1}{a} \right) \quad (14)$$

“vis viva” equation (“life force”)

discovered prior to concept of energy

handy: gives total speed  $v$  at any radius  $r$

*Q: at which  $r$  is  $v = 0$ ? how does this work for a circular orbit?*

*Q: for a given orbit (fixed  $e$ ), when is  $v$  max?*

## Unbound Orbits

Note that both parabolic and hyperbolic orbits are not periodic – do not close on themselves  
“one-way ticket” past the central object

### Parabola

$$e = 1$$

$$r = \frac{2p}{1 + \cos \theta} \quad (15)$$

$p$  is distance of closest approach

for parabolic orbit:

$TE = 0$  exactly!  $\rightarrow KE = -PE$  exactly! very special case!

$$\Rightarrow GM/r = \frac{1}{2}v^2$$

So at  $r = \infty$ ,  $v = 0$

to have this orbit, launch from  $r$  with speed

$$v_{\text{launch}} = \sqrt{2GM/r}$$

## iClicker Poll: Orbits

given: test particle  $m$ , at distance  $r$  from gravitating body  $M$   
for test particle to have total energy  $TE = 0$   
launch from  $r$  with speed  $v_0 = \sqrt{2GM/r}$

Q: what happens if launch with speed  $v > v_0$ ?

- A particle will be in a bound orbit: circle or ellipse
- B particle will be unbound, with speed  $v \rightarrow 0$  as  $r \rightarrow \infty$
- C particle will be unbound, with speed  $v > 0$  as  $r \rightarrow \infty$

Q: *why is  $v_0$  a special speed?*

## Escape Speed

At radius  $r$ , define

**escape speed**

$$v_{\text{esc}} = \sqrt{2GM/r}$$

- if launch from  $r$  with  $v_{\text{launch}} < v_{\text{esc}}$   
then  $TE < 0$ : fall back! (elliptical orbit)
- if launch from  $r$  with  $v_{\text{launch}} > v_{\text{esc}}$   
then  $TE > 0$ : escape “easily”:  $v > 0$  at  $r = \infty$
- if launch from  $r$  with  $v_{\text{launch}} = v_{\text{esc}}$  exactly  
then  $TE = 0$  exactly, “just barely” escape

So: escape speed is *minimum speed* needed to leave a gravitating source

Example: escape speed from earth

$$v_{\text{esc}} = 11 \text{ km/s} = 25,000 \text{ mph!}$$

predict the future: if toss object with  $v < 25,000$  mph, falls back  
but if  $v > 25,000$  mph  $Q$ : *example?* never returns!

finally, the more “generic” unbound orbit:

## hyperbola

$$r(\theta) = \frac{(e^2 - 1)a}{1 + e \cos \theta} \quad (16)$$

$$e > 1, TE > 0$$

$v > 0$  at  $r = \infty$ : nonzero speed far from  $M$

Recall: at large  $r$ , hyperbola  $\rightarrow$  *straight line*

But Newton says:  $d\vec{v}/dt = -GM/r^2 \hat{r}$

so as  $r \rightarrow \infty$ , then  $d\vec{v}/dt \rightarrow 0$

$\Rightarrow$  gravity negligible,  $\vec{v} \rightarrow$  const: *free body*=straight line!

orbit of unbound “flyby”:

- $\nabla$  distant nearly free body  $\rightarrow$  passing: pulled toward  $M$   
 $\rightarrow$  distant deflected nearly free body