

# Astro 501: Radiative Processes

## Lecture 14

Feb 15, 2013

Announcements:

- **Problem Set 4** today at 5pm
- **Problem Set 5** due next Friday

Last time: Thomson scattering

*Q: what is Thomson scattering?*

*Q: what does the scattered power depend on?*

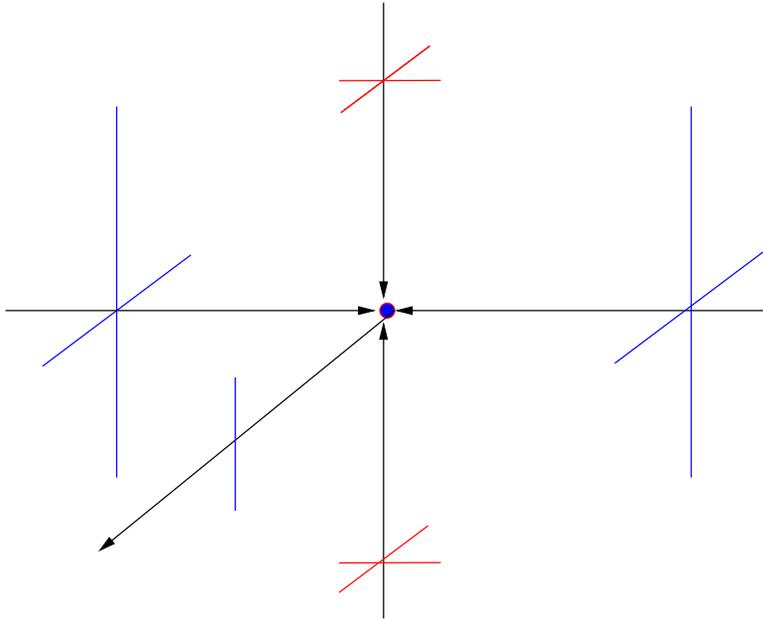
*Q: what does  $d\sigma/d\Omega$  depend on? and not?*

*Q: what does  $\sigma$  depend on? and not?*

*Q: lessons?*

*Q: when does Thomson scattering generate polarization?*

if initial radiation field has quadrupole intensity pattern



linear polarization!

lesson: polarization arises from Thomson scattering  
when electrons “see” quadrupole anisotropies in radiation field

# Awesomest Example of Thompson Polarization: the CMB

The CMB is nearly isotropic radiation field  
arises from  $\tau = 1$  “surface of last scattering” at  $z = 1000$   
when free  $e$  and protons “re” combined  $ep \rightarrow H$

- *before recombination:*

  - Thomson scattering of CMB photons, Universe opaque

- *after recombination:* no free  $e$ , Universe transparent

consider electron during last scatterings  
sees and anisotropic thermal radiation field

consider point at hot/cold “wall”

  - locally sees *dipole*  $T$  anisotropy

  - net polarization towards us: zero! Q: *why?*

ω

Q: *what about edge of circular hot spot? cold spot?*

polarization tangential (ring) around hot spots  
radial (spokes) around cold spots  
(superpose to “+” = zero net polarization—check!)

www: WMAP polarization observations of hot and cold spots

Note: polarization &  $T$  anisotropies *linked*

→ consistency test for CMB theory and hence hot big bang

## Polarization Observed

First detection: pre-WMAP!

★ DASI (2002) ground-based interferometer  
at level predicted based on  $T$  anisotropies! Woo hoo!

WMAP (2003): first polarization- $T$  correlation function

Planck (March 2013): much more sensitive to polarization  
maybe a signature of inflation-generated gravitational radiation?

# Bremsstrahlung

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German lesson for today:

*Bremse* = break (as in stopping)

*Strahlung* = radiation

→ Bremsstrahlung = “breaking radiation”  
= radiation from decelerated charge particles

Consider a **dilute plasma** at temperature  $T$ , with

- free ions: charge  $+Ze$ , number density  $n_i$
- free electrons: charge  $-e$ , number density  $n_e$

Q: *astrophysical examples?* www: awesome example

Q: *what microphysics what will cause the plasma to emit?*

∨ *i.e., what interactions will occur?*

Q: *which particles will radiate more?*

dilute plasma = low particle density = typical in astrophysics  
→ three-body collisions unlikely; ignore these  
→ focus on two-body collisions

possible interactions: Coulomb forces between particle pairs

- electron-electron
- ion-ion
- electron-ion

But note: for *two identical charged particles*

dipole moment in center of mass  $\vec{d} = \sum q_i \vec{r}_i = 0$

*no dipole radiation*

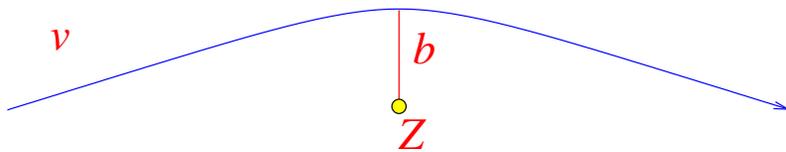
So: *electron-ion* dipole radiation dominates

$\infty$  electron and ion scattered by same Coulomb force (Newton III)  
But  $a_i/a_e = m_e/m_i < 10^{-3}$  → ion acceleration negligible  
→ focus electron acceleration in static field of ion

# Order of Magnitude Expectations

start with *classical, nonrelativistic* picture

consider a free, unbound electron with asymptotic speed  $v$  moving in Coulomb field of stationary ion



let  $b =$  *the distance of closest approach* or **impact parameter**

Q: *estimate of maximum acceleration?*

Q: *duration of acceleration? velocity change? radiation frequency*

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Recall the *Spirit of Order-of-Magnitude*:

- ignore all dimensionless constants, e.g., “small circle approximation”  $2\pi \approx 1$
- lower expectations for precision
- use rough result to guide more careful calculations

maximum acceleration: Coulomb acceleration at closest approach

$$a_{\max} \sim \frac{Ze^2}{m_e b^2} \quad (1)$$

duration of acceleration: **collision time**

$$\tau \sim \frac{b}{v} \quad (2)$$

velocity change

$$\Delta v \sim a_{\max} \tau \sim \frac{Ze^2}{m_e b v} \sim \left( \frac{Ze^2/b}{m_e v^2} \right) v \quad (3)$$

frequency of radiation: use only timescale in problem

$$\omega \sim \frac{1}{\tau} \sim \frac{v}{b} \quad (4)$$

10 Q: what is maximum radiated power? radiated energy? energy per unit freq?

maximum radiated power is

$$P_{\max} \sim \frac{e^2 a_{\max}^2}{c^3} \sim \frac{e^2 \Delta v^2}{c^3 \tau^2} \sim \frac{Z^2 e^6}{m_e v^2 b^2 \tau^2} \quad (5)$$

radiated energy

$$\Delta W \sim P_{\max} \tau \sim \frac{Z^2 e^6}{m_e v^2 b^2 \tau} \quad (6)$$

radiated energy per unit frequency

$$\frac{\Delta W}{\Delta \nu} \sim \frac{\Delta W}{\omega} \sim \frac{Z^2 e^6}{m_e v^2 b^2} \quad (7)$$

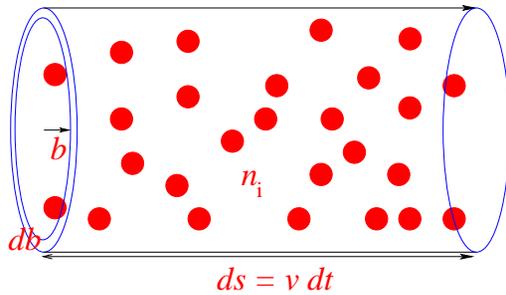
this energy radiated per electron-ion encounter at distance  $b$

electron with speed  $v$  moves encounters ion number density  $n_i$

- we want number of ions  $dN_i$  that  $e$  encounters

≡ out to distance  $\sim b$  in time  $dt$  *Q: which is?*

- *Q: what is typical rate of energy emitted per electron?*



in cylindrical distance  $(b, b + db)$ , volume swept is

$$dV = 2\pi b db ds = 2\pi v b db dt \quad (8)$$

i.e.,  $dV \sim b^2 v dt$

thus number of ions encountered is

$$d\mathcal{N}_i = n_i dV \sim n_i b^2 v dt \quad (9)$$

Thus the rate of energy emitted = *power emitted per e* is

$$\frac{dP_{\text{pere}}}{d\nu} = \frac{\Delta W}{\Delta\nu} \frac{d\mathcal{N}_i}{dt} \sim \frac{e^6 Z^2}{m_e c^3 v} n_i \quad (10)$$

Q: and so what is emission coefficient  $j_\nu$ ?

Our order-of-magnitude estimate for the emission coefficient from nonrelativistic bremsstrahlung:

$$j_\nu = n_e \frac{dP_{\text{pere}}}{d\nu} \sim \frac{e^6 Z^2}{m_e c^3 \nu} n_e n_i \quad (11)$$

*Q: what's the basic physical picture?*

*Q: notable features? what didn't we get from order of mag?*

*Q: how can we do the classical calculation more carefully?*