

Astro 501: Radiative Processes

Lecture 21

March 6, 2013

Announcements:

- **Problem Set 6** due Friday 5pm
- Midterm Exam: grading elves hard at work

Last time: synchrotron spectrum

Q: spectrum for isotropic, monoenergetic electrons?

Q: spectrum for electrons with a power-law energy distribution?

define **critical frequency**

$$\omega_c \equiv \frac{3}{2}\gamma^3\omega_B \sin \alpha = \frac{3}{2}\gamma^2 \frac{qB \sin \alpha}{mc} = \frac{3}{2}\gamma^2 \omega_g \sin \alpha \quad (1)$$

$$\nu_c = \frac{\omega_c}{2\pi} = \frac{3}{4\pi}\gamma^3\omega_B \sin \alpha \quad (2)$$

emission spectrum is synchrotron function $F(\omega/\omega_c)$

sharply peaked near $\omega_c \propto \omega_g \gamma^2$

full expression for power-law electron spectrum
of the form $dN/d\gamma = C\gamma^{-p}$

$$4\pi j_{\text{tot}}(\omega) = \frac{\sqrt{3}q^3 C B \sin \alpha}{2(p+1)\pi m c^2} \Gamma\left(\frac{p}{4} + \frac{9}{12}\right) \Gamma\left(\frac{p}{4} - \frac{1}{12}\right) \left(\frac{mc\omega}{3qB \sin \alpha}\right)^{-(p-1)/2} \quad (3)$$

with $\Gamma(x)$ the gamma function, with $\Gamma(x+1) = x \Gamma(x)$

Q: *expected spectral index?*

Q: *do you expect the signal to be polarized? how?*

Polarization of Synchrotron Radiation

for an electron with a single pitch angle $\tan \alpha = v_{\perp}/v_{\parallel}$

→ circular motion around field line

→ radiation circularly polarized orthogonal to \vec{B}
and elliptically polarized at arbitrary angles

but with distribution of pitch angles α ,

elliptical portion cancels out → partial **linear polarization**

polarization strength varies with projected angle
of magnetic field on sky

more power orthogonal to projected field direction

→ net linear polarization, detailed formulae in RL

averaging over power-law distribution of electron energies

↳ partial polarization is $\Pi = (p + 1)/(p + 7/3)$

and so $\Pi = 3/4$ for $p = 3$: highly polarized!

Transition from Cyclotron to Synchrotron

How and why are the emission spectra so different for cyclotron (non-relativistic) vs synchrotron (relativistic)?

recall: in either case, electron motion is *strictly periodic* with angular frequency

$$\omega_B = \frac{qB \sin \alpha}{m c \gamma} \quad (4)$$

Q: *nature of Fourier spectrum of received field?*

Q: *Fourier spectrum of emission for single pitch angle?*

Q: *spectrum in nonrelativistic case $\gamma \rightarrow 1$?*

Q: *spectrum in mildly relativistic case?*

electron motion at fixed α strictly periodic with ω_B
→ received field also strictly periodic
→ Fourier transform of field is nonzero only for
discrete *series* of frequencies $m\omega_B$, $m \in 1, 2, \dots$

and thus received radiation also is a Fourier series in ω_B

cyclotron = nonrelativistic case: see field $E = E_0 \cos \omega_B t$
Fourier series has *one term*: the fundamental frequency ω_B

when mildly relativistic: Doppler effects add harmonic at $2\omega_B$
and electric field shape modified to sharper, narrower peak

going to strongly relativistic: many harmonics excited
series “envelope” approaches $F(\omega/\omega_c)$
electric field → very sharp, very narrow peak

○ with distribution of pitch angles:
“spaces” in series filled in → continuous spectrum

Synchrotron Self-Absorption

Recall strategy so far:

- calculate emission coefficient j_ν
- remember Kirchoff's law $j_\nu = \alpha_\nu B_\nu(T)$
- solve for $\alpha_\nu = j_\nu / B_\nu(T)$

We have already found

Q: why won't this work here?

Q: what do we need to do? hint—how did we handle a two-level system?

Kirchoff's law is only good for a *thermal* system where emitter and absorber particles are nonrelativistic and have Maxwell-Boltzmann energy/momentum distribution

here: electrons are relativistic and nonthermal

really: Kirchoff is example of *detailed balance*

→ in equilibrium, emission and absorption rates are the same → this still applies in nonthermal case

recall from *2-level* system, with $E_2 = E_1 + h\nu$

$$\alpha_\nu \stackrel{\text{2-level}}{=} \frac{h\nu}{4\pi} [n(E_1)B_{12} - n(E_2)B_{21}] \phi(\nu) \quad (5)$$

Q: *physical interpretation of $n(E_1)$? B_{12} ? B_{21} ? $\phi(\nu)$?*

∞

Q: *how should this be modified for synchrotron electrons?*

in *2-level system*, emission at frequency ν
arises from *unique energy level spacing* $E_2 = E_1 + h\nu$

but cosmic ray electrons have *continuous energy spectrum*
→ emission at ν can arise from *any two energies*:
generalized to

$$\alpha_\nu = \frac{h\nu}{4\pi} \sum_{E_1} \sum_{E_2} [n(E_1)B_{12} - n(E_2)B_{21}] \phi_{21}(\nu) \quad (6)$$

- with $\phi_{21}(\nu) \rightarrow \delta[\nu - (E_2 - E_1)/h]$
- first term: true absorption
- second term: stimulated emission

the goal: recast this in terms of what we know

⊙ synchrotron emission j_ν

we have

$$\alpha_\nu = \frac{h\nu}{4\pi} \sum_{E_1} \sum_{E_2} [n(E_1)B_{12} - n(E_2)B_{21}] \phi_{21}(\nu) \quad (7)$$

use Einstein relations, good for thermal and nonthermal

- spontaneous emission rate from state E_2 : $A_{21} = 2h\nu^3 B_{21}/c^2$
- absorption and stimulated emission: $B_{21} = B_{12}$

note that spontaneous *emission* is what we know!

we have found synchrotron power $P(\nu, E_2) = 2\pi P(\omega)$,
with E_2 the radiating electron's energy

$$P(\nu, E_2) = h\nu \sum_{E_2} A_{21} \phi_{21}(\nu) \quad (8)$$

now impose Einstein conditions and simplify

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Q: role of ϕ_{21} and double sum $\sum_{E_1} \sum_{E_2}$?

profile function $\phi_{21}(\nu) \rightarrow \delta(E_2 - E_1 - h\nu)$
 fixes E_1 for a given E_2 and ν
 and double sum \rightarrow single sum

$$\alpha_\nu = \frac{c^2}{8\pi h\nu^3} \sum_{E_2} [n(E_2 - h\nu) - n(E_2)] P(\nu, E_2) \quad (9)$$

so far: schematic sum over electron energies
 but really a continuum

recall: in each phase space cell h^3

- number of electron states with momentum p is $g_e f(p)$
 - volume density of states in momentum space volume is d^3p/h^3
- and thus

$$\alpha_\nu = g_e \frac{c^2}{8\pi h\nu^3} \frac{1}{h^3} \int [f(p_2^*) - f(p_2)] P(\nu, E_2) d^3p_2 \quad (10)$$

where p_2^* is the momentum corresponding to energy $E_2 - h\nu$

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Q: how is f related to electron spectrum $N(E)$?

number of electrons per unit volume
with energy in $(E, E + dE)$ is $N(E) dE$

but this means that

$$N(E) dE = \frac{4\pi g_e}{h^3} p^2 f(p) dp \quad (11)$$

and for ultrarelativistic electrons, $E = cp$

thus we have

$$\alpha_\nu = \frac{c^2}{8\pi h\nu^3} \int \left[\frac{N(E - h\nu)}{(E - h\nu)^2} - \frac{N(E)}{E^2} \right] E^2 P(\nu, E) dE \quad (12)$$

and since $h\nu \ll E$, expand to first order

$$\alpha_\nu = -\frac{c^2}{8\pi\nu^2} \int dE P(\nu, E) E^2 \partial_E \left[\frac{N(E)}{E^2} \right] \quad (13)$$

and for a power-law $N(E) \propto E^{-p}$, we have

$$-E^2 \partial_E \left[\frac{N(E)}{E^2} \right] = (p + 2) \frac{N(E)}{E} \quad (14)$$

Synchrotron Absorption

finally then

$$\alpha_\nu = (p + 2) \frac{c^2}{8\pi\nu^2} \int dE P(\nu, E) \frac{N(E)}{E} \quad (15)$$

note frequency dependence:

- prefactor ν^{-2}
 - integral $\int dE P(\nu)N(E)/E \sim dE P(\nu)E^{-(p+1)} \sim \nu^{-p/2}$
- net scaling: $\alpha_\nu \propto \nu^{-(p+4)/2}$

full result

$$\alpha_\nu = \frac{\sqrt{3}}{8\pi} \Gamma\left(\frac{3p+2}{12}\right) \Gamma\left(\frac{3p+22}{12}\right) \left(\frac{3q}{2\pi m^3 c^5}\right)^{p/2} \left(\frac{q^3 C}{m}\right) (B \sin \alpha)^{(p+2)/2} \nu^{-(p+4)/2}$$

Source Function

source function

$$S_\nu = \frac{j_\nu}{\alpha_\nu} \propto \frac{\nu^{-(p-1)/2}}{\nu^{-(p+4)/2}} = \nu^{5/2} \quad (16)$$

to see this, recall that

$$j_\nu \sim \int dE N(E) P(\nu) \quad (17)$$

$$\alpha_\nu \sim \nu^{-2} \int dE \frac{N(E)}{E} P(\nu) \quad (18)$$

thus source function has

$$S_\nu \sim \nu^2 \bar{E} \quad (19)$$

with typical electron energy $\bar{E} = m\bar{\gamma}$ for freq ν

but $\nu(E) \approx \nu_c(E) \sim E^2$, so $\bar{E} \propto \nu^{1/2}$

and thus $S_\nu \sim \nu^{5/2}$ independent of electron spectral index

Synchrotron Radiation: the Big Picture

for relativistic electrons with power-law energy distribution

emission coefficient

$$j_\nu \propto \nu^{-(p-1)/2} \quad (20)$$

absorption coefficient

$$\alpha_\nu \propto \nu^{-(p+4)/2} \quad (21)$$

source function (note nonthermal character!)

$$S_\nu \propto \nu^{5/2} \quad (22)$$

Q: optical depth vs ν ? implications?

Q: spectrum of a synchrotron emitter?

www: awesome example: pulsar wind nebulae

young pulsars are spinning down

much of rotational energy goes into relativistic wind

which collides with the supernova ejecta and emits synchrotron