

Astro 501: Radiative Processes
Lecture 25
March 15, 2013

Announcements:

- **Problem Set 8** due today
- **Problem Set 7** due after break

Last time: Sunyaev-Zeldovich theory

Q: implications of Planck form of CMB?

Q: what is the SZ effect? integrated intensity?

Q: effect at low frequencies?

CMB has Planck (blackbody) form $I_\nu = B_\nu(T_0)$

recall: a blackbody spectrum arises from

- a *thermal emitter* having source function $S_\nu = B_\nu$
- that is also *optically thick*

thus we conclude: sometime in the past

- cosmic matter and radiation were *in thermal equilibrium*
- and the Universe was *opaque*

but the fact that the CMB is a *background*

to low- z objects \rightarrow late-time U. is *transparent* to CMB

thus the CMB implies that the Universe is *evolving*

and in the past was much *denser*

so that equilibrium could be established

SZ: Low Frequencies

at low frequencies $h\nu \ll kT_e$, we have

$$\frac{\Delta n(\nu)}{n_0(\nu)} = \frac{\Delta I_\nu}{I_\nu^0} \approx -2y \quad (1)$$

- *frequency-independent fractional decrease* in intensity
- proportional to Compton y

physically reasonable? yes!

these wimpy photons are promoted to higher frequencies

Note: low-frequencies are Rayleigh-Jeans regime
using antenna temperature $T_\nu = c^2 I_\nu / 2hk\nu^2$, we have

$$\frac{\Delta T_\nu}{T_\nu} \approx -2y \quad (2)$$

ω constant fractional decrease in antenna temperature

Q: what about the high-frequency limit $h\nu \gg kT_e \sim m_e c^2 \beta^2$?

At *high frequencies* $h\nu \gg kT_e \sim \beta^2$, *Compton shift dominates*

roughly expect upscattered frequency

$$\nu \approx \left(1 + \frac{h\nu_0}{m_e c^2}\right) \nu_0 \quad (3)$$

so that $\Delta\nu/\nu_0 \sim h\nu_0/m_e c^2$

same song, second verse:

scattering kernel, again in simpleminded approximation

$$K(\nu, \nu_0) = (1 - \tau) \delta(\nu_0 - \nu) + \tau \delta\left[\nu_0 - \left(1 - h\nu/m_e c^2\right)\nu\right]$$

thus we have

$$\frac{\Delta n(\nu)}{n(\nu)} = \frac{\Delta I_\nu}{I_\nu} = -\tau \frac{h\nu}{m_e c^2} \frac{1}{n(\nu)} \nu \partial_\nu n(\nu) \quad (4)$$

$$= -y \frac{h\nu}{kT_{\text{rad}}} \left(2 - \frac{h\nu/kT_{\text{rad}} e^{h\nu/kT_{\text{rad}}}}{e^{h\nu/kT_{\text{rad}}} - 1}\right) \approx y \left(\frac{h\nu}{kT_{\text{rad}}}\right)^2 \quad (5)$$

Q: *implications?*

SZ Effect: High Frequencies

for $h\nu \gg kT_{\text{rad}}$, a careful calculation finds

$$\frac{\Delta I_\nu}{I_\nu} \approx y \left(\frac{h\nu}{kT_{\text{rad}}} \right)^2 \quad (6)$$

now *fractional shift is frequency dependent*

To summarize, we have found the limiting SZ shifts for $x = h\nu/kT_{\text{rad}}$, we have

$$\frac{\Delta I_\nu}{I_\nu} \rightarrow \begin{cases} -2y & x \ll 1 \\ yx & x \gg 1 \end{cases} \quad (7)$$

Q: *implication?*

SZ Effect: Null

We see that SZ *decreases low-frequency intensity* and *increases high-frequency intensity*

→ there must be a transition that crosses zero!

→ frequency ν_{null} must exist at which $\Delta I_\nu / I_\nu = 0$: **SZ null**

PS7 (modulo typos): Kompaneets SZ approximation

$$\frac{\Delta I_\nu}{I_\nu} = -y \, g(x) = -y \frac{x e^x}{e^x - 1} \left[4 - x \frac{e^x + 1}{e^x - 1} \right] \quad (8)$$

with $x = h\nu/kT_{\text{rad}}$; perturbation vanishes at $x_{\text{null}} = 3.83$

Q: SZ spectral dependence?

Q: how to exploit this observationally?

o www: SZ spectrum plotted

Q: effect of cosmic redshifting on SZ: integrated intensity? spectrum?

SZ and Cosmic Redshifting

redshifting: a fundamental consequence the expanding universe cosmologist learn at their mother's knee that *lengths* stretch with to the *cosmic scale factor* $a(t)$:

$$\vec{r}(t) = a(t) \vec{r}_0 \quad (9)$$

usual convention: today (time t_0) distance is \vec{r}_0 , so $a(t_0) = 1$

wavelengths are lengths – it's right there in the name!

so wavelengths stretch as $\lambda(t) \propto a(t)$

- photon emitted with λ_{em} at t_{em} is observed at $t_{\text{obs}} = t_0$ with

$$\lambda_{\text{obs}} = \frac{a(t_{\text{obs}})}{a(t_{\text{em}})} \lambda_{\text{em}} = \frac{1}{a(t_{\text{em}})} \lambda_{\text{em}} \quad (10)$$

- thus the *redshift* is related to scale factor via

$$z \equiv \frac{\lambda_{\text{obs}}}{\lambda_{\text{em}}} - 1 = \frac{1}{a(t_{\text{em}})} - 1 \quad (11)$$

Q: redshifting effect on ν ? T ? SZ $\Delta I_\nu / I_\nu$? $\Delta I / I$?

Cosmic Redshifting

wavelength: $\lambda_{\text{obs}} = \lambda_{\text{em}}/a_{\text{em}} = (1 + z)\lambda_{\text{em}}$

frequency: $\nu = c/\lambda$ and so $\nu_{\text{obs}} = a_{\text{em}} \nu_{\text{em}} = \nu_{\text{em}}/(1 + z)$

temperature: Wein says $T/\nu_{\text{max}} = \text{const}$

so $T_{\text{obs}} \propto \nu_{\text{max,obs}} \propto a_{\text{em}}$ and so $T_{\text{obs}} = T_{\text{em}}/(1 + z)$

But recall SZ fractional intensity change $\Delta I_{\nu}/I_{\nu}^0 = -yg(x)$

→ only depends on dimensionless ratio x evaluated at emission

$$x = \frac{h\nu_{\text{em}}}{kT_{\text{rad,em}}} = \frac{h\nu_{\text{obs}}/(1 + z)}{kT_{\text{rad,obs}}/(1 + z)} = \frac{h\nu_{\text{obs}}}{kT_{\text{rad,obs}}} \quad (12)$$

SZ fractional specific intensity change is *redshift independent!*

∞ → same goes for integrated intensity change $\Delta I/I = 4y$

Q: *is this reasonable?*

www: SZ clusters at different redshift

this also means:

- shape of spectral distortion is redshift independent
- max, min, and null in $\Delta I_\nu / I_\nu$ always the same!
- robust signature of and test of (thermal) SZ effect!

SZ clusters usually resolved

Q: what information is in a cluster's SZ pattern on sky?

Q: what information is in total SZ flux across sky?

Q: SZ applications for cluster astrophysics?

Thermal SZ Effect as a Probe of Galaxy Cluster

in each line of sight

SZ measures Comptonization parameter in a cluster:

$$y = \sigma_T \int \frac{n_e kT_e}{m_e c^2} ds = \frac{\sigma_T}{m_e c^2} \int P_e ds \approx \frac{\sigma_T kT_e}{m_e c^2} \int n_e ds \quad (13)$$

direct measurement of *projected pressure* in column
and if T_e known, a measure of electron column density

SZ flux measures

$$\int \cos \theta y d\Omega \approx \int y d\Omega = \frac{\int y dA}{D_A^2} \quad (14)$$

where $D_A(z)$ is the (angular diameter) distance

$$\int y dA \approx \frac{\sigma_T kT_e}{m_e c^2} \int n_e ds dA \propto M_{\text{gas}} \quad (15)$$

→ SZ flux gives *intracluster cluster gas mass!* Q: cosmo apps?

SZ Effect: Cosmological Applications

- *SZ identifies all clusters without redshift bias!*
→ SZ can be used to discover high- z clusters
- SZ + X-ray gives cluster size, gas mass, T_e
if cluster physics well-understood (Ricker, Vijayaraghavan)
→ *cluster mass*
- cluster number density (“abundance”) and mass vs z
i.e., cluster *mass function* a sensitive probe of cosmology

today: clusters are the *largest bound objects*; in early U: rare number and mass vs time sensitive to *cosmic acceleration* that competes with *structure growth via gravitational instability*
⇒ clusters probe this competition

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Q: so how to find clusters, measure redshifts?

note that SZ redshift independence also means
SZ does not give cluster redshift

Dark Energy Survey key project:
optical images, redshifts of clusters
compare with SZ survey by South Pole Telescope

www: SPT survey image

SZ Effect: More Cosmological Applications

even for clusters not clearly imaged in SZ
SZ effect from all clusters still imprinted on CMB
affects ΔT_{cmb} perturbation pattern on sky

typical angular size of cluster SZ:

for large cluster $\theta_{\text{cluster}} \sim R_{\text{cluster}}/d_{\text{H}} \sim 3 \text{ Mpc}/4 \text{ Gpc} \sim 3 \text{ arcmin}$

i.e., SZ affects small angular scales

in C_ℓ multipole space this corresponds to $\ell \sim 200/\theta_{\text{deg}} \sim 4000$

SZ statistical imprint on CMB anisotropies:

exquisitely sensitive measure of *cosmic structure*

for experts: angular power spectrum $C_\ell^{\text{SZ}} \propto \sigma_8^7!$

To date: SZ contribution to power spectrum not seen! *Planck?*

Kinetic SZ Effect and Cosmology

Thus far: implicitly assumed that cluster is *at rest* relative to CMB frame (“fundamental observers”)

but if cluster moving along line of sight with velocity v_{los} bulk motion adds *uniform Doppler shift*

to usual thermal SZ effect

→ kinematic or **kinetic SZ effect** (“kSZ”)

at low frequencies (Rayleigh-Jeans), kSZ has

$$\frac{\Delta I_\nu}{I_\nu} \approx -\tau \frac{v_{\text{los}}}{c} \quad (16)$$

what causes cluster motion = bulk flows?

→ large-scale density perturbations

Q: but how do we tell between kSZ and thermal SZ?

Other SZ Applications

SZ is sensitive to *any* population of high-energy electrons

should be a SZ contribution from relativistic electrons in intracluster medium, arising from

- cosmic-rays escaped from galaxies, or from
- “structure formation cosmic rays” accelerated by cosmological shocks

nonthermal SZ

also recall that AGN jets lead to electron acceleration
in principle can measure jet electrons in SZ

Have a good break!