

Astro 501: Radiative Processes

Lecture 10

Feb 6, 2013

Announcements:

- **Problem Set 3** due next time
- BDF office hours shortened today
- TA office hours 1:30-3:00 tomorrow

Last time: began classical EM radiation

Q: energy density?

Q: Poynting vector?

└ Today: plane waves & polarization

Maxwell and Fourier Modes

We have seen: wave equation demands $\omega = ck$
But Maxwell equations impose further constraints

Consider arbitrary Fourier modes

$$\vec{E} = E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} \hat{a}_1, \text{ and } \vec{B} = B_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} \hat{a}_2$$

Maxwell equations in vacuum impose conditions:
for example, Coulomb's law $\nabla \cdot \vec{E} = 0$ implies

$$\vec{k} \cdot \vec{E} = 0 \tag{1}$$

or equivalently $\hat{n} \cdot \hat{a}_1 = 0$

similarly, no monopoles requires

$$\vec{k} \cdot \vec{B} = 0 \quad \hat{n} \cdot \hat{a}_2 = 0 \tag{2}$$

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Q: what does this mean physically for the waves?

we found $\vec{k} \cdot \vec{E} = \vec{k} \cdot \vec{B} = 0$

→ propagation orthogonal to field vectors

⇒ *EM waves are transverse*

Faraday's law requires $\omega \vec{B} = c \vec{k} \times \vec{E}$, or

$$\vec{B} = \frac{c \vec{k}}{\omega} \times \vec{E} = \hat{n} \times \vec{E} \quad (3)$$

and Ampère's law gives $\vec{E} = -\hat{n} \times \vec{B}$

Q: what do these conditions imply for the waves?

Faraday's law gives $\vec{B} = \hat{n} \times \vec{E}$, so

$$\vec{E} \cdot \vec{B} = \vec{E} \cdot (\hat{n} \times \vec{E}) = 0 \quad (4)$$

$\Rightarrow \vec{E}$ and \vec{B} are orthogonal to each other!

Faraday also implies

$$|B|^2 = \hat{n}^2 |E|^2 - |\hat{n} \cdot \vec{E}|^2 = |E|^2 \quad (5)$$

using vector identity $(\hat{a} \times \vec{b}) \cdot (\vec{c} \times \vec{d}) = \hat{a} \cdot \vec{c} \vec{b} \cdot \vec{d} - \hat{a} \cdot \vec{d} \vec{b} \cdot \vec{c}$

we have: $E_0 = B_0$: field amplitudes are equal

which in turn means: $\hat{a}_2 = \hat{n} \times \hat{a}_1$, and $\hat{a}_1 \cdot \hat{a}_2 = 0$

→ $(\hat{n}, \hat{a}_1, \hat{a}_2)$ form an orthogonal basis

Monochromatic Plane Wave: Time Averaging

at a given point in space, field amplitudes vary sinusoidally with time \rightarrow energy density and flux also sinusoidal but we are interested in timescales $\gg \omega^{-1}$:
 \rightarrow take *time averages*

Useful to use *complex* field amplitudes
then take *real part* to get physical component

handy theorem: for $A(t) = \mathcal{A}e^{i\omega t}$ and $B(t) = \mathcal{B}e^{i\omega t}$
i.e., same time dependence, then time-averaged products

$$\langle \text{Re}A(t) \text{Re}B(t) \rangle = \frac{1}{2}\text{Re}(\mathcal{A}\mathcal{B}^*) = \frac{1}{2}\text{Re}(\mathcal{A} * \mathcal{B}) \quad (6)$$

Monochromatic Plane Wave: Energy, Flux

time-averaged Poynting flux amplitude

$$\langle S \rangle = \frac{c}{8\pi} \text{Re}(E_0 B_0^*) = \frac{c}{8\pi} |E_0|^2 = \frac{c}{8\pi} |B_0|^2 \quad (7)$$

time-averaged energy density

$$\langle u \rangle = \frac{c}{8\pi} |E_0|^2 = \frac{c}{8\pi} |B_0|^2 \quad (8)$$

Q: given wave direction \vec{n} , degrees of freedom in \vec{E}, \vec{B} ?

Polarization

EM waves propagating *in a particular direction* \hat{n}
must be transverse $\vec{k} \cdot \vec{E} = \hat{n} \cdot \vec{E} = 0$
→ nonzero \vec{E} components lie in plane \perp to \hat{n}

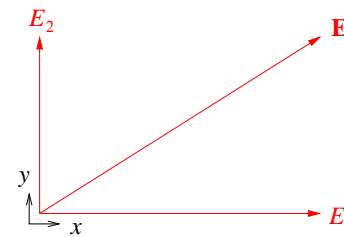
two independent components

for convenience: choose coordinates where $\hat{n} = \hat{z}$
so fields are in transverse plane $x - y$
physical electric vector is *real part* of

$$\vec{E} = (E_1 \hat{x} + E_2 \hat{y}) e^{-i\omega t} \quad (9)$$

complex amplitudes can be written

$$E_1 = \mathcal{E}_1 e^{i\phi_1} \quad E_2 = \mathcal{E}_2 e^{i\phi_2} \quad (10)$$



Q: *but wait—what about the magnetic field?*

transverse electric field has two independent components
but once \vec{E} determined, then $\vec{B} = \hat{n} \times \vec{E}$
at every point along sightline \hat{n} , magnetic \perp electric
 \Rightarrow *no additional degrees of freedom for \vec{B}*

monochromatic plane wave *has two independent components*

consider plane at fixed $z = \hat{n} \cdot \vec{r}$, say $z = 0$
the two *physical* components of the field evolve as

$$E_x = \mathcal{E}_1 \cos(\omega t - \phi_1) \quad E_y = \mathcal{E}_2 \cos(\omega t - \phi_2) \quad (11)$$

with E_1, E_2 can take any values, and ϕ_1, ϕ_2 independent
but only difference $\phi_1 - \phi_2$ can be important
 \rightarrow a total of *3 independent parameters* describe the fields

- ∞ Q: \vec{E} time evolution if E_1 and E_2 can differ, but $\phi_1 - \phi_2 = 0$?
Q: same but $\phi_1 - \phi_2 = \pi$?

Linear Polarization

For $\phi_1 - \phi_2 = 0$, we have

$$E_x = \mathcal{E}_1 \cos(\omega t - \phi_1) \quad E_y = \mathcal{E}_2 \cos(\omega t - \phi_1) = \frac{\mathcal{E}_2}{\mathcal{E}_1} E_x \quad (12)$$

fields share same sign and same sinusoidal time dependence

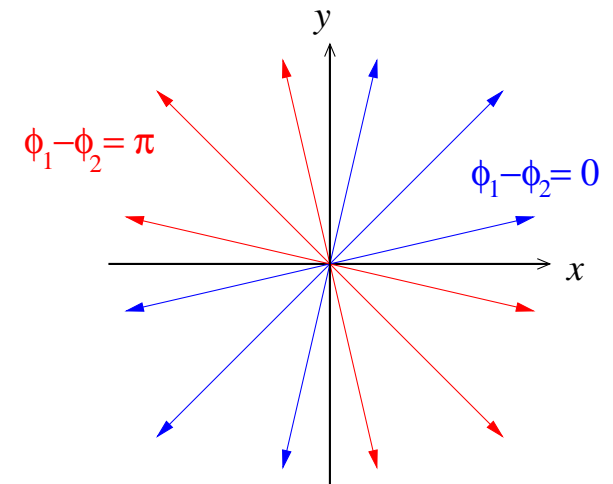
\vec{E} sweeps out *line with positive slope* in $x - y$ plane

→ **linear polarization**

For $\phi_1 - \phi_2 = \pi$, fields share time dependence

but have **opposite sign**

→ linear polarization with negative slope



Q: what is \vec{E} time dependence if

$\mathcal{E}_1 = \mathcal{E}_2$ but $\phi_1 - \phi_2 = \pi/2$ or $-\pi/2$

Circular Polarization

if $\mathcal{E}_1 = \mathcal{E}_2$ but $\phi_1 - \phi_2 = \pi/2$

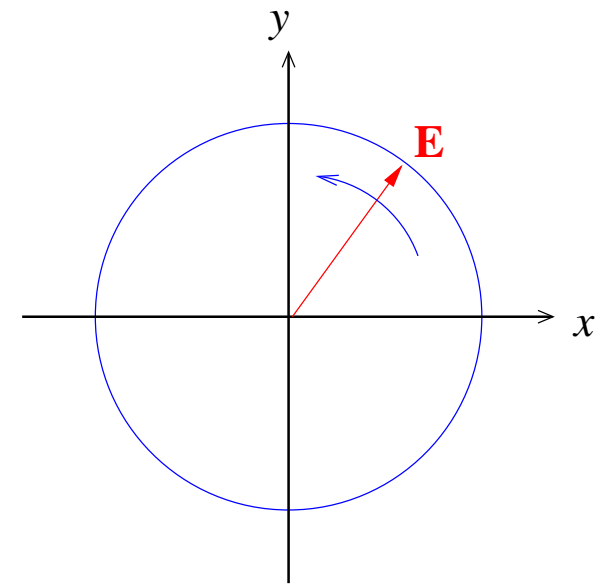
$$E_x = \mathcal{E}_1 \cos(\omega t - \phi_1) \quad E_y = \mathcal{E}_1 \sin(\omega t - \phi_1)$$

\vec{E} sweeps counterclockwise circle
as seen approaching observer

⇒ **circular polarization**

Engineering: “*lefthanded*” circular polarization

→ but using righthand rule: *positive helicity*



if $\mathcal{E}_1 = \mathcal{E}_2$ but $\phi_1 - \phi_2 = -\pi/2$

→ “*righthand*” circular polarization, or *negative helicity*

in the most general case: $\mathcal{E}_1 \neq \mathcal{E}_2$ and $\phi_1 - \phi_2$ arbitrary

Q: what is \vec{E} time dependence?

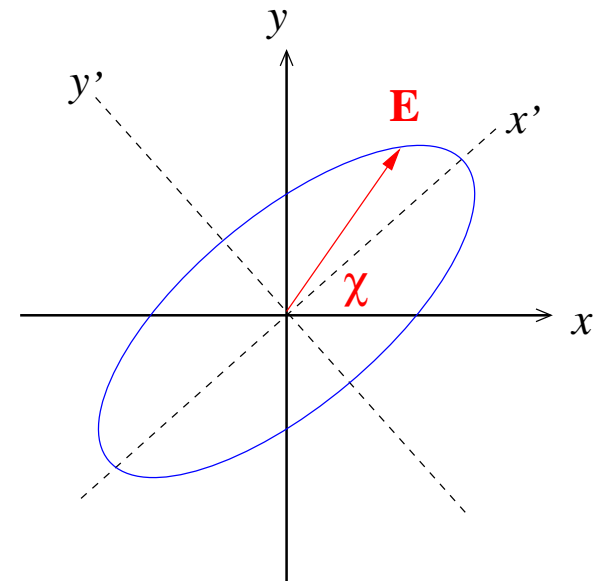
Elliptical Polarization

in the general case

$$E_x = \mathcal{E}_1 \cos(\omega t - \phi_1) \quad E_y = \mathcal{E}_2 \cos(\omega t - \phi_2)$$

intuitively, blends linear and circular features:

→ **elliptical polarization**



ellipse *orientation* fixed by $\mathcal{E}_1 - \mathcal{E}_2$ difference

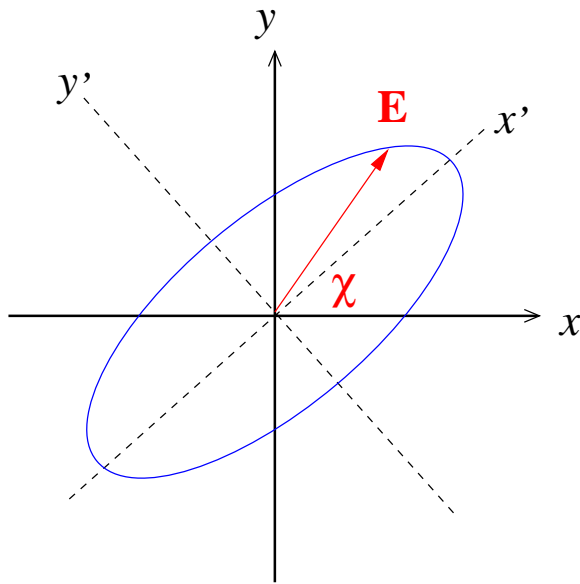
ellipse *eccentricity* and *helicity* fixed by $\phi_1 - \phi_2$ difference

in coordinates (x', y') rotated to align with *principal axes*

$$E'_x = \mathcal{E}_0 \cos \beta \cos(\omega t) \quad E'_y = \mathcal{E}_0 \sin \beta \sin(\omega t)$$

11 for some $\beta \in [-\pi/2, +\pi/2]$

Q: evolution if $\beta > 0$?



$$E'_x = \mathcal{E}_0 \cos \beta \cos(\omega t) \quad E'_y = -\mathcal{E}_0 \sin \beta \sin(\omega t)$$

principle axes: $\mathcal{E}_0 \cos \beta$ and $\mathcal{E}_0 \sin \beta$

if $\beta \in [0, \pi/2]$: ellipse sweeps clockwise

→ “*righthanded*” elliptical polarization, *negative helicity*

if $\beta \in [\pi/2, \pi]$: “*lefthanded*”, *positive helicity*

Q: what give linear polarization? circular?

we want to relate $x - y$ field parameters

$\mathcal{E}_1, \mathcal{E}_2, \phi_1, \phi_2$

with principle axes $x' - y'$ parameters $\mathcal{E}_0, \beta, \chi$

rotate $x - y$ components by angle χ

$$E_x = \mathcal{E}_0 (\cos \beta \cos \chi \cos \omega t - \sin \beta \sin \chi \sin \omega t)$$

$$E_y = \mathcal{E}_0 (\cos \beta \sin \chi \cos \omega t - \sin \beta \cos \chi \sin \omega t)$$

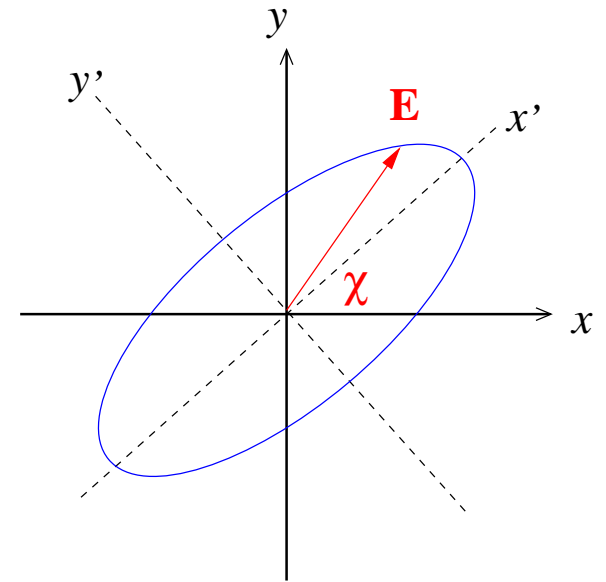
matching to, e.g., $E_x = \mathcal{E}_1 \cos(\omega t - \phi_1)$:

$$\mathcal{E}_1 \cos \phi_1 = \mathcal{E}_0 \cos \beta \cos \chi \quad (13)$$

$$\mathcal{E}_1 \sin \phi_1 = \mathcal{E}_0 \sin \beta \sin \chi \quad (14)$$

$$\mathcal{E}_2 \cos \phi_2 = \mathcal{E}_0 \cos \beta \sin \chi \quad (15)$$

$$\mathcal{E}_2 \sin \phi_2 = -\mathcal{E}_0 \sin \beta \cos \chi \quad (16)$$



Q: how can we determine polarization by intensity measurements?

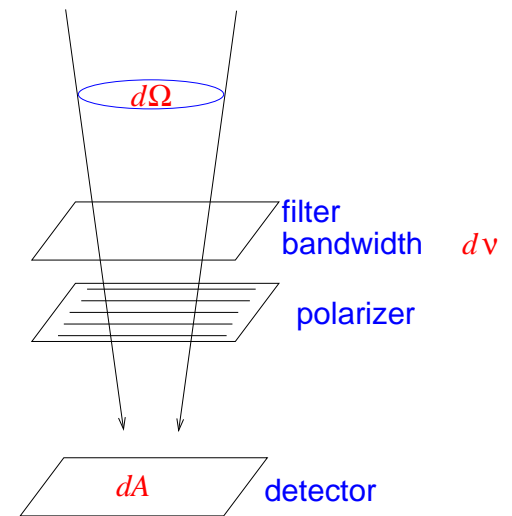
Introduce *polarizer*

can *rotate* polarizer:

→ measure I_x , I_y , and 45° rotated $I_{x'}$, $I_{y'}$

can use circular polarizers to measure

→ positive and negative circular polarization I_+ , I_-



combine: **Stokes parameters**

$$I = I_x + I_y \quad (17)$$

$$Q = I_x - I_y \quad (18)$$

$$U = I_{x'} - I_{y'} \quad (19)$$

$$V = I_+ - I_- \quad (20)$$

14 Q: what physically is each? can more than one of Q, U, V be nonzero? what does that correspond to?

Q: range of values for Q ? U ? V ? are they all independent?

Stokes Parameters

for *monochromatic waves*, Stokes parameters related to $\mathcal{E}_1, \mathcal{E}_2, \phi_1, \phi_2$ and $\mathcal{E}_0, \beta, \chi$ bases:

$$I = \mathcal{E}_1^2 + \mathcal{E}_2^2 = \mathcal{E}_0^2 \quad (21)$$

$$Q = \mathcal{E}_1^2 - \mathcal{E}_2^2 = \mathcal{E}_0^2 \cos 2\beta \cos 2\chi \quad (22)$$

$$U = 2\mathcal{E}_1\mathcal{E}_2 \cos(\phi_1 - \phi_2) = \mathcal{E}_0^2 \cos 2\beta \sin 2\chi \quad (23)$$

$$V = 2\mathcal{E}_1\mathcal{E}_2 \sin(\phi_1 - \phi_2) = \mathcal{E}_0^2 \sin 2\beta \quad (24)$$

and thus

$$\mathcal{E}_0 = \sqrt{I} \quad (25)$$

$$\sin 2\beta = V/I \quad (26)$$

$$\tan 2\chi = U/Q \quad (27)$$

since wave has 3 independent parameters,
Stokes parameters must be *related*

$$I^2 = Q^2 + U^2 + V^2 \quad (28)$$

Quasi-Monochromatic Waves

natural light generally *not a pure monochromatic wave*
with a single, definite, complete state of polarization

rather: a *superposition* of components with many polarizations

consider wave with *slowly varying* amplitudes and phases

$$E_1(t) = \mathcal{E}_1(t) e^{i\phi_1(t)} ; \quad E_2(t) = \mathcal{E}_2(t) e^{i\phi_2(t)} \quad (29)$$

“slow”: wave looks completely polarized on timescale ω^{-1}
but amplitudes and phases drift over intervals $\Delta t \gg \omega^{-1}$
→ polarization changes

but also wave is *no longer monochromatic*

frequency spread: “*bandwidth*” $\Delta\omega \sim 1/\Delta t \ll \omega$

→ *quasi-monochromatic wave*

Q: effect on Stokes?

Stokes Parameters for Quasi-Monochromatic Light

real measurements represent **averages** over timescales during which polarization can change

Stokes parameters become averages

$$I = \langle E_1 E_1^* \rangle + \langle E_2 E_2^* \rangle = \langle \mathcal{E}_1^2 + \mathcal{E}_2^2 \rangle \quad (30)$$

$$Q = \langle E_1 E_1^* \rangle - \langle E_2 E_2^* \rangle = \langle \mathcal{E}_1^2 - \mathcal{E}_2^2 \rangle \quad (31)$$

$$U = \langle E_1 E_2^* \rangle + \langle E_2 E_1^* \rangle = 2\langle \mathcal{E}_1 \mathcal{E}_2 \cos(\phi_1 - \phi_2) \rangle \quad (32)$$

$$V = -i (\langle E_1 E_2^* \rangle - \langle E_2 E_1^* \rangle) = 2\langle \mathcal{E}_1 \mathcal{E}_2 \sin(\phi_1 - \phi_2) \rangle \quad (33)$$

but for quasi-monochromatic waves

$$I^2 \geq Q^2 + U^2 + V^2 \quad (34)$$

- quasi-monochromatic polarization is still in general **elliptical**
- but drifts can reduce degree of polarization

$$I^2 \geq Q^2 + U^2 + V^2 \quad (35)$$

- maximum polarization when equality holds: *completely elliptically polarized*
- minimum when $Q = U = V = 0$: *unpolarized*
- arbitrary wave is *partially polarized*

useful to define *polarized* intensity

$$I_{\text{pol}} = Q^2 + U^2 + V^2 \quad (36)$$

and since $I_{\text{pol}} \leq I$, define fractional **degree of polarization**

$$\Pi \equiv \frac{I_{\text{pol}}}{I} = \frac{\sqrt{Q^2 + U^2 + V^2}}{I} \quad (37)$$

note: can always decompose Stokes parameters

$$\begin{pmatrix} I \\ Q \\ U \\ V \end{pmatrix} = \begin{pmatrix} I - I_{\text{pol}} \\ 0 \\ 0 \\ 0 \end{pmatrix} + \begin{pmatrix} I_{\text{pol}} \\ Q \\ U \\ V \end{pmatrix} \quad (38)$$

sum of unpolarized and polarized components

Superposition and Stokes

consider composite wave that is superposition of many independent waves

electric field components are given by **superposition**

$$E_1 = \sum_k E_1^{(k)} \quad ; \quad E_2 = \sum_k E_2^{(k)} \quad (39)$$

each term k of which has different phase

PS3: phases specified, can calculate sum explicitly

but generally, *phases are random*

so field products average out phases from different waves

$$\langle E_i E_j^* \rangle = \sum_k \sum_\ell \langle E_i^{(k)} E_j^{(\ell)*} \rangle = \sum_k \langle E_i^{(k)} E_i^{(k)*} \rangle \quad (40)$$

but due to this averaging, *Stokes parameters are additive*

$$I = \sum_k I^{(k)} \quad (41)$$

$$Q = \sum_k Q^{(k)} \quad (42)$$

$$U = \sum_k U^{(k)} \quad (43)$$

$$V = \sum_k V^{(k)} \quad (44)$$