

Astro 507
Lecture 22
March 12, 2014

Announcements:

- **Preflight 4 due 9am Friday**

Last time:

recombination \rightarrow huge drop in free $e^- \rightarrow$ CMB freeze/decouple
to calculate in detail: need cosmic statistical mechanics

key inputs: uncertainty principle, Boltzmann factor,

and baryon-to-photon ratio $\eta = n_b/n_\gamma$

Q: *whatsa baryon? η order of magnitude?*

key outputs: non-rel, non-degen $n = g(mT/2\pi\hbar^2)^{3/2} e^{-(m-\mu)/T}$

for reaction in (“chemical”) equilibrium: $\sum \mu_i = \sum \mu_f$

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Q: *apply to recombination?*

The Mighty Saha Equation

Recombination: equal forward and reverse rates for



and so chem potentials have

$$\mu_p + \mu_e = \mu_{\text{H}} \quad (1)$$

for non-rel species $n = g(mT/2\pi\hbar^2)^{3/2}e^{-(m-\mu)/T}$

thus we have **Saha equation**

$$\frac{n_e n_p}{n_{\text{H}}} = \frac{g_e g_p}{g_{\text{H}}} \left(\frac{m_e m_p}{m_{\text{H}}} \right)^{3/2} \left(\frac{T}{2\pi\hbar^2} \right)^{3/2} e^{-(m_e + m_p - m_{\text{H}})/T} \quad (2)$$

$$\approx \left(\frac{m_e T}{2\pi\hbar^2} \right)^{3/2} e^{-B/T} \quad (3)$$

where $B \equiv m_e + m_p - m_{\text{H}} = 13.6 \text{ eV}$

introduce “free electron fraction” $X_e = n_e/n_B$

use $n_B = \eta n_\gamma \propto \eta T^3$

from Extras last time: $n_\gamma = 2\zeta(3)/\pi^2 T^3$, with $\zeta(3) = \sum_1^\infty 1/n^3 = 1.20206\dots$

and note that $n_p = n_e$ Q: *why?*, so

$$\frac{n_e^2}{n_H n_B} = \frac{X_e^2}{1 - X_e} = \frac{\sqrt{\pi}}{4\sqrt{2}\zeta(3)\eta} \left(\frac{m_e}{T}\right)^{3/2} e^{-B/T} \quad (4)$$

Q: *sanity checks? what sets characteristic T scale?*

Q: *when is $X_e = 0$ (exactly)?*

At last–recombination!

Q: *how define physically?*

Q: *how define operationally, in terms of X_e ?*

ω Q: *given some $X_{e,\text{rec}}$, how to get z_{rec} ?*

The Epoch of Recombination

Saha gives

$$\frac{1 - X_e}{X_e^2} = \frac{4\sqrt{2}\zeta(3)}{\pi^{1/2}} \eta \left(\frac{B}{m_e}\right)^{3/2} \left(\frac{T}{B}\right)^{3/2} e^{B/T} \quad (5)$$

if always equilib, then strictly $X_e = 0$ only at $T = 0$
but note $e^{B/T}$: X_e exponentially small when $T \ll B$

viewed as a function of $B/T \equiv u$

$$\frac{1 - X_e}{X_e^2} = \frac{4\sqrt{2}\zeta(3)}{\pi^{1/2}} \eta \left(\frac{B}{m_e}\right)^{3/2} u^{3/2} e^u \equiv A u^{3/2} e^u \quad (6)$$

where $A = 4\sqrt{2}/\pi^{1/2}\zeta(3) \eta (B/m_e)^{3/2}$

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- Q: *what is order-of-magnitude of A?*
 - Q: *implications for recombination?*
 - Q: *physical picture?*

in recombination Saha expression $(1 - X_e)/X_e = A(B/T)^{3/2}e^{B/T}$
prefactor is tiny!

$$A \sim \eta(B/m_e)^{3/2} \sim 10^{-9}(10^{-5})^{3/2} \sim 10^{-16} !$$

why? largely due to *tiny baryon-to-photon ratio*

but when recombine: $1 - X_e \simeq X_e$

so require $1 \sim 10^{-16}(B/T_{\text{rec}})^{3/2}e^{B/T_{\text{rec}}}$

\Rightarrow so need $B/T_{\text{rec}} \gg 1$ to offset prefactor

\Rightarrow and thus $T_{\text{rec}} \ll B$!

more carefully define recomb: $X_e = X_{e,\text{rec}} = 0.1$

(arbitrary, but not crazy; see PS4)

then solve for T_{rec} :

$$\frac{B}{T_{\text{rec}}} = \ln \left(\frac{\pi^{1/2}}{4\sqrt{2}\zeta(3)} \right) + \ln \left(\frac{1 - X_{e,\text{rec}}}{X_{e,\text{rec}}^2} \right) + \ln \eta^{-1} + \frac{3}{2} \ln \frac{m_e}{B} + \frac{3}{2} \ln \frac{B}{T}$$

5 $\sim 40 \quad (\gg 1)$

(ignore or iterate $\ln B/T$ term)

Recombination Quantified

and so

$$T_{\text{rec}} \approx \frac{B}{40} \simeq 0.3 \text{ eV} \ll B \quad (7)$$

$$z_{\text{rec}} \approx 1400 \ll z_{\text{rec,naive}} \quad (8)$$

$$t_{\text{rec}} \approx \frac{2}{3\sqrt{\Omega_m}} H_0^{-1} (1 + z_{\text{rec}})^{-3/2} = 350,000 \text{ yrs} \quad (9)$$

PS4: try it yourself!

Implications for CMB frequency spectrum:

- at recomb: emission lines created at $h\nu_{\text{rec}} \gtrsim 3B/4$
and thus at $h\nu_{\text{rec}} \gtrsim 30kT_{\text{rec}}$
- post-recomb: T and ν both redshift the same way, so
- CMB spectrum *distorted* from Planck at high freq: $h\nu \gtrsim 30kT$
- small signal, difficult to observe, but tantalizing www: predictions

Q: what physically is responsible for $T_{\text{rec}} \ll B$?

Recombination “Delay”

Why is $T_{\text{rec}} \ll B$?

- ▷ because for small X_e , Saha says $X_e \propto 1/\eta^{1/2} \gg 1$
- ▷ many photons per baryon: even if typically $E_\gamma \ll B$, high-E tail of Planck distribution not negligible (at first)
lots of **ionizing photons** with $E_\gamma \geq B$
H dissociated as soon as formed

When does dissociation stop?

can show that fraction of photons with $E_\gamma > B$

is roughly $f_{\text{ionizing}} \sim e^{-B/T}$

so ratio of **ionizing** photons per baryon is

$$\frac{n_{\gamma,\text{ionizing}}}{n_B} \sim \frac{e^{-B/T}}{\eta} \quad (10)$$

estimate recombination when $n_{\gamma,\text{ionizing}}/n_B \sim 1$

→ $T \sim B/\ln \eta^{-1} \ll B$ (check!)

⇒ **recombination “delayed”** to huge photon-to-baryon ratio

Recombination: Hydrogen Level Population

recall: Boltzmann expression for atomic hydrogen (H I):

$$\frac{n(2P)}{n(1S)} = 3e^{-3B_H/4T} = 3e^{-120,000 \text{ K}/T} \quad (11)$$

Q: implications for H populations?

consider recombining $p + e \rightarrow H + \gamma$ throughout recomb:

Q: what is γ energy at emission?

Q: what happens to γ ?

Q: implications?

Recombination: Nonequilibrium Effects

for $p + e \rightarrow \text{H}(n = 1) + \gamma$:

- $E_\gamma = B_{\text{H}}$ “Lyman limit”
 - H atoms *absorption cross section huge* at this energy
photon mean free path $\ell = 1/n_{\text{H}}\sigma_{\text{abs}}$ tiny
universe optically thick to Lyman photons
- \Rightarrow quickly reionizes another H atom! *no net change!*

To overcome delay

- recombine to 1st excited state: $p + e \rightarrow \text{H}(2p) + \gamma$
- single photon $\text{H}(2p) \rightarrow \text{H}(1s) + \gamma$ Ly α transition
also optically thick, also no net progress
- *two-photon transition* $\text{H}(2p) \rightarrow \text{H}(1s) + \gamma + \gamma$ can go
but probability & rate smaller than for single photon
- eventually redshifting takes Lyman photons off resonance

net effect: *delays recombination relative to Saha*

Last Scattering: Photons Decouple from Matter

“recombination” a smooth transition in X_e , not instantaneous

www: equilibrium X_e plot

nevertheless, exponential drop in X_e around z_{rec}

photons interact with gas via Thomson scattering: $\gamma e \rightarrow \gamma e$

rate per photon of scattering with e :

$$\Gamma_e(\gamma) = n_e \sigma v = n_e \sigma_T c = X_e n_b \sigma_T c \quad (12)$$

drop in $X_e \rightarrow$ abrupt slowdown in scattering

as usual, competition between interaction and expansion
interactions “stop” when

$$\Gamma_e(\gamma) \lesssim H \quad (13)$$

and solving for $\Gamma_e(T) = H(T)$ gives last scattering :

$$z_{\text{ls}} \sim 1100 \quad (14)$$

After last scattering:

- photons “decoupled” from gas
- but $X_e \neq 0$: some free e, p remain

Q: what is X_e as $T \rightarrow 0$? why?

Freezing of Recombination

when typical photon has last scattering with e
still some residual ionization: i.e., some free e, p
can they recombine? yes!

do they recombine? yes, for a short while...then no!

Why? recombination rate per p : $\Gamma_{\text{rec},p} \sim n_e \sigma_{\text{rec}} v_{\text{therm}}$
with $\sigma_{\text{rec}} \sim (m_e/T) \sigma_{\text{T}}$ and $v_{\text{therm}} \sim \sqrt{T/m_e}$
recombination stops when $\Gamma_{\text{rec},p} \lesssim H$

after this: cooling does not reduce ionization

fixed value of $X_{e,\text{freeze}} \sim 10^{-4}$: “freeze-in of residual ionization”
at

$$z_{\text{ri}} \simeq 1000 \quad (15)$$

Q: cosmological implications of $X_{e,\text{freeze}} \neq 0$?

Recombination Timeline Summarized

The large drop in free electron density around $z \sim 1000$ leads to three distinct but related events:

(1) recombination U. **ionized** \rightarrow **neutral**

$$X_e \rightarrow X_{e,\text{rec}} \sim 0.1: z_{\text{rec}} \sim 1300$$

...but photons still coupled to gas, and vice versa

(2) last scattering typical photons no longer interacts with e

U. **opaque** \rightarrow **transparent**

$$\Gamma_e(\gamma) \sim H: z_{\text{ls}} \sim 1100$$

...but gas still coupled to photons Q : *how can this be?*

$$T_{\text{gas}} = T_{e,p,H} = T_\gamma$$

(3) residual ionization freeze-in

free e and p diluted until “can’t find each other”

But even still: photons scatter off residual ionization
 e and thus p, H still exchange energy
with thermal photon bath: $T_{e,p,H} = T_\gamma$ still!
when does this stop?

(4) gas decoupling

typical residual e no longer has photon interactions
gas decouples from photons

when? Thomson scattering rate *per e* : $\Gamma_e = n_\gamma \sigma_T c \lesssim H$
at $z_{\text{dec,gas}} \sim 500$

note: scatter rate *per e* $\Gamma_e \gg \Gamma_\gamma$ = *scatter rate per CMB photon*

Summary of CMB Highlights

CMB Observed

can make precision observations of spectrum, sky distribution thanks to sophisticated radio techniques and instruments

- CMB fantastically isotropic: $\delta T/T \sim \text{few} \times 10^{-5}$
- CMB exquisitely thermal

CMB Theory

detailed, precise calculations of recomb, last scattering, thanks well-known atomic physics

- isotropic CMB \rightarrow U. was once very homogeneous
- Planckian CMB spectrum \rightarrow U. was once thermalized
 \rightarrow plasma hot, dense enough to equilibrate

CMB \rightarrow demands hot big bang in FLRW universe!

Extrapolated current U to $t \sim 400,000$ yr

and $z \sim 1000 \rightarrow$ **great success!**

Emboldens us to push earlier!

Primordial Nucleosynthesis

Prelude to Nucleosynthesis

Q: what sets T scale for element (nuclei) synthesis?

Q: what component dominates cosmic density, expansion then?

Q: what is the particle content of the universe then?

Q: what form(s) do the baryons take then? mesons?

Nucleosynthesis: Setting the Stage

★ light elements formed in nuclear reactions
relevant scale: nuclear binding energies \sim MeV

★ $T \sim$ MeV at redshift $z_{\text{bbn}} = T/T_0 - 1 \sim 10^{10}$!
since $z_{\text{bbn}} \gg z_{\text{eq}}$ (matter-rad equality)
well into radiation dominated era: $\rho \approx \rho_{\text{rad}}$

www: Ω vs a plot

will see: $t(1 \text{ MeV}) \sim 1 \text{ sec}$

★ particle content at BBN

relativistic species: photons, neutrinos, e^\pm when $T \gtrsim m_e$

non-relativistic species: baryons, e^- when $T \ll m_e$

what about dark matter? energy?

DM presumably non-rel, weakly interacting: inert during BBN

DE: also assume not important for dynamics, microphysics

...but can later relax these assumptions and test them!

Who Feels What? Particles and Forces

$$\begin{pmatrix} u \\ d \\ e \\ \nu_e \end{pmatrix} \begin{pmatrix} c \\ s \\ \mu \\ \nu_\mu \end{pmatrix} \begin{matrix} \text{charm quark} \\ \text{strange quark} \\ \text{mu lepton (muon)} \end{matrix} \begin{pmatrix} t \\ b \\ \tau \\ \nu_\tau \end{pmatrix} \begin{matrix} \text{top quark} \\ \text{bottom quark} \\ \text{tau lepton} \end{matrix} \quad (16)$$

quarks: feel all fundamental forces (strong, EM, weak, gravity)
 carry conserved quantum number: **baryon number**

leptons: do *not* feel strong force

but also carry conserved quantum number: **lepton number**

▷ *charged* leptons: feel EM, weak, gravity

▷ neutrinos: only feel weak, gravity

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More bragging rights:

in BBN, *all four* fundamental forces play a crucial role!