

Astro 210
Lecture 8
Feb. 2, 2018

Announcements

- **HW2 due online in PDF, Today 5:00 pm**
- HW3 available, due next Friday
- HW1 grading under way; you'll be notified when done
- **register** your iClicker; link on course webpage
- first Planetarium shows Mon Feb 5 and Wed Feb 7
info online: **reservations**, schedules, directions, report form

Last Time: Sir Isaac Weighs In

Newton's Laws of motion

I. inertia *Q: just a special case of Newton II?*

II. $\vec{F} = m\vec{a}$ *Q: fortunetelling & archæology?*

III. action-reaction

Q: when/where/to what do Newton's laws of motion apply?

Newtonian Universal Gravitation

*Q: magnitude of gravity force between masses m and M
with distance r ?*

Q: direction of the force?

Q: why universal?

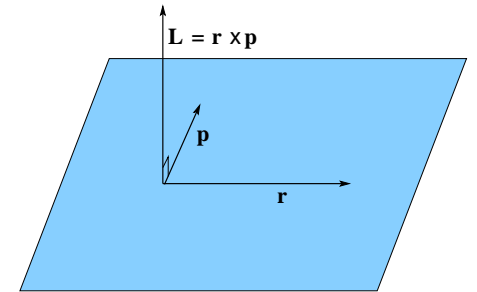
Q: what's a variable? what's a constant?

Angular Momentum

For point mass, **angular momentum** defined as:

$$\vec{L} = \vec{r} \times \vec{p} = m\vec{r} \times \vec{v} \quad (1)$$

i.e., using cross product



look at time change:

$$\frac{d}{dt}\vec{L} = m\dot{\vec{r}} \times \vec{v} + m\vec{r} \times \dot{\vec{v}} \quad (2)$$

$$= m\vec{v} \times \vec{v} + m\vec{r} \times \vec{a} \quad (3)$$

$$= \vec{r} \times \vec{F} = \vec{\tau} \quad \text{torque} \quad (4)$$

angular counterpart of Newton II:

- net (linear) force changes linear momentum
- net twisting force = torque changes angular momentum

Gravity and Angular Momentum

angular momentum changed by net torque

$$\frac{d}{dt}\vec{L} = \vec{r} \times \vec{F} \quad (5)$$

when force is due to **gravity**, torque:

$$\vec{\tau} = \vec{r} \times \vec{F} = -G\frac{mM}{r^3}\vec{r} \times \vec{r} = 0 \quad (6)$$

so if force is gravity, then

$$\frac{d}{dt}\vec{L} = 0 \quad (7)$$

and thus $\vec{L} = \text{const}$:

↳ angular momentum is **conserved!**

Q: *what about gravity force guaranteed this?*

What Keeps the Earth in Orbit?

consider *circular orbit* → requires centripetal acceleration
angular speed $d\theta/dt = \omega = 2\pi/P = \text{const}$

$$\vec{a}_c = -\omega^2 \vec{r} = -\frac{v^2}{r} \hat{r}$$

diagram: show \vec{v} , \vec{r} , \vec{a}

Newton II: acceleration demands net force

Newton gravity supplies a force!

→ Newtonian gravity is crucial and necessary ingredient
for understanding the dynamics of planetary motion
but have to see how the detailed predictions
compare with observation

Testing Newton with Kepler

Program:

- assume Newtonian gravity controls planetary motion
- that is, for any planet let $\vec{F}_{\text{net}} = \vec{F}_{\text{Sun-planet}}$
- input this into Newton's laws
- turn mathematical cranks \rightarrow predict orbits
- compare predictions with observations *Q: how?*

Solutions: Orbits

For attractive inverse square force law, Newton II orbits are cross sections of cone:

- circle
- ellipse
- parabola
- hyperbola

Circle eccentricity $e = 0$

at each point:

$$F = ma = mv_C^2/r$$

$$\Rightarrow GMm/r^2 = mv_C^2/r$$

\Rightarrow circular orbits have speed

$$v_C = \sqrt{\frac{GM}{r}}$$

\checkmark example: find circular speed 1 AU from Sun

$$v_C = 3 \times 10^4 \text{ m/s}$$

Kepler from Newton

Kepler I: Orbits are ellipses

with Sun at one focus

Newton II + Universal Inverse-Square Gravity

bound orbits around a gravitational source M are

- ellipses
- with M at one focus
- eccentricity e set by angular momentum L

Check!

Kepler from Newton

Kepler II: Equal areas in equal times

Newton: consider small time interval dt

move angle $d\theta = \omega dt$

sweep area

diagram: top view: path, $d\theta$, \vec{r} , \vec{v} , \vec{v}_θ

$$dA = \frac{1}{2}r^2 d\theta = \frac{1}{2}r^2 \omega dt \quad (8)$$

but $\omega = v_\theta/r$, where $\vec{v}_\theta \perp \vec{r}$

\Rightarrow swept area

$$dA = \frac{1}{2}r^2 \frac{v_\theta}{r} dt = \frac{1}{2}rv_\theta dt \quad (9)$$

⇒ swept area

$$dA = \frac{1}{2} r^2 \frac{v_\theta}{r} dt = \frac{1}{2} r v_\theta dt \quad (10)$$

finally, $r v_\theta = |\vec{r} \times \vec{v}| = |\vec{L}|/m$

Q: *why?*, so

$$dA = \frac{1}{2} \frac{L}{m} dt \quad (11)$$

Woo hoo! were' home free! Q: *why?*

But $L = \text{const}$ for radial force ($\vec{r} \times \vec{F} = 0$)

so

$$\frac{dA}{dt} = \frac{L}{2m} = \text{const} \quad (12)$$

Kepler II!

Check!

And we learn:

Kepler II comes from angular momentum conservation!

Note: in absence of other forces

and for perfectly spherical Sun

ellipses close on themselves—path repeated forever

Kepler III: $a^3 = kP^2$

Newton: can prove generally for elliptical orbits
bad news: price is lotsa algebra

good news: simple to do for circular orbits
circular $\rightarrow r = a$, and $v^2 = GM/a$
but also $v = 2\pi a/P$ Q: why?

$$v^2 = \left(\frac{2\pi a}{P}\right)^2 = \frac{4\pi^2 a^2}{P^2} \quad (13)$$

$$= \frac{GM}{a} \quad (14)$$

$$\Rightarrow a^3 = \left(\frac{GM}{4\pi^2}\right) P^2 \quad (15)$$

check!

¹² bonus: $k = GM/4\pi^2$ depends on mass of central object
 \rightarrow same k for all planets

Energy

For “test” particle m moving due to gravity of M
Gravitational potential energy: Q : why “potential”?

$$PE = -GMm/r$$

Kinetic energy:

$$KE = \frac{1}{2}m\vec{v}^2 = \frac{1}{2}m(v_x^2 + v_y^2 + v_z^2) = \frac{1}{2}m\dot{r}^2 \quad (16)$$

Total energy:

$$TE = KE + PE = -GMm/r + \frac{1}{2}m\dot{r}^2$$

key result: $d(TE)/dt = 0$

→ total energy conserved!

that is: value of TE the same for all time!

Orbits Revisited

Bound orbits (circle & ellipse): in polar coordinates

$$r(\theta) = \frac{(1 - e^2)a}{1 + e \cos \theta} \quad (17)$$

Circle radius $r = a = \text{const}$, eccentricity $e = 0$

recall: circular orbit has constant speed $v_C^2 = GM/r$

$$PE = -\frac{GMm}{r} < 0 \quad (18)$$

$$KE = \frac{1}{2}mv_C^2 = \frac{1}{2}m\frac{GM}{r} = \frac{1}{2}\frac{GMm}{r} = -\frac{1}{2}PE \quad (19)$$

$$\Rightarrow TE = KE + PE = PE/2 = -|PE|/2 < 0 \quad (20)$$

$TE < 0$: negative? yes!

Q: what does it mean to have negative energy?

for orbiting system $TE < 0$:

→ have to *supply* energy to system to break it apart

Why? when particles are *at rest* and “very” *far apart*

$$KE = mv^2/2 = 0$$

$$PE = GMm/r \rightarrow 0 \quad Q: \text{how far apart is this?}$$

and so $TE = KE + PE = 0$: zero total energy

But if start in closed orbits (circular or elliptical): $TE < 0$

→ To “break” the system from closed orbits, must *add* energy

But energy is conserved → not spontaneously added

so system is **bound**

⇒ can't fall apart without external influence

Note: $KE = -PE/2 = |PE|/2$ generally true for

gravitating systems in equilibrium:

“virial theorem”

ellipse: semimajor axis a , eccentricity $0 < e < 1$

turns out: TE depends only on a , not e

from conservation of energy

$$TE = -GMm/r + \frac{1}{2}mv^2 = -GMm/2a < 0 \rightarrow \text{bound}$$

can show

$$v^2 = GM \left(\frac{2}{r} - \frac{1}{a} \right) \quad (21)$$

“vis viva” equation (“life force”)

discovered prior to concept of energy

handy: gives total speed v at any radius r

Q: at which r is $v = 0$? how does this work for a circular orbit?

Q: for a given orbit (fixed e), when is v max?

Unbound Orbits

Note that both parabolic and hyperbolic orbits are not periodic – do not close on themselves
“one-way ticket” past the central object

Parabola

$$e = 1$$

$$r = \frac{2p}{1 + \cos \theta} \quad (22)$$

p is distance of closest approach

for parabolic orbit:

$TE = 0$ exactly! $\rightarrow KE = -PE$ exactly! very special case!

$$\Rightarrow GM/r = \frac{1}{2}v^2$$

So at $r = \infty$, $v = 0$

to have this orbit, launch from r with speed

$$v_{\text{launch}} = \sqrt{2GM/r}$$

iClicker Poll: Orbits

given: test particle m , at distance r from gravitating body M
for test particle to have total energy $TE = 0$
launch from r with speed $v_0 = \sqrt{2GM/r}$

Q: what happens if launch with speed $v > v_0$?

- A particle will be in a bound orbit: circle or ellipse
- B particle will be unbound, with speed $v \rightarrow 0$ as $r \rightarrow \infty$
- C particle will be unbound, with speed $v > 0$ as $r \rightarrow \infty$

Q: *why is v_0 a special speed?*

Escape Speed

At radius r , define

escape speed

$$v_{\text{esc}} = \sqrt{2GM/r}$$

- if launch from r with $v_{\text{launch}} < v_{\text{esc}}$
then $TE < 0$: fall back! (elliptical orbit)
- if launch from r with $v_{\text{launch}} > v_{\text{esc}}$
then $TE > 0$: escape “easily”: $v > 0$ at $r = \infty$
- if launch from r with $v_{\text{launch}} = v_{\text{esc}}$ exactly
then $TE = 0$ exactly, “just barely” escape

So: escape speed is *minimum speed* needed to leave a gravitating source

Example: escape speed from earth

$$v_{\text{esc}} = 11 \text{ km/s} = 25,000 \text{ mph!}$$

predict the future: if toss object with $v < 25,000$ mph, falls back
but if $v > 25,000$ mph Q : *example?* never returns!