Astro 596/496 NPA Lecture 18 February 27, 2019

Announcements:

Problem Set 3 out, due next Friday
 Office Hours Wed 2-3pm or by appointment
 if questions arise will post Discussion on Compass

Last time: neutrons and protons in the early Universe *Q: when in equilibrium, what is n/p ratio? Q: what reactions maintain this equilibrium? Q: why does this equilibrium end? when?* 

### **Problem Set 3: Hints and Errata**

- for numerical answers,  $\eta = 6 imes 10^{-10}$  but you can use  $\eta \sim 10^{-9}$  without penalty
- Question 2(b): can skip last part involving  $\Omega_B h^2$
- Cosmic Entropy

note that energy density is  $\varepsilon = \varepsilon(T) = (\partial E / \partial_V)_T$ so that in 2nd law of thermo,  $E = E(T, V) = \varepsilon(T) V$ similarly S = s(T) V, N = n(T) V

in notes, factor of 2 typo fixed:

Planck function for photon number density distribution

$$\frac{dn}{dE} = \frac{g}{2\pi^2 \hbar^3 c^2} E^2 f(E) = \frac{g}{2\pi^2 \hbar^3 c^2} \frac{E^2}{e^{E/kT} - 1}$$
(1)

Ν

in early universe:  $T \gg 1$  MeV:

ω

neutrino and electron interactions allow nucleon interconversion

$$n + \nu_e \leftrightarrow p + e^-$$
 (2)

$$p + \bar{\nu}_e \leftrightarrow n + e^+$$
 (3)

finds equilibrium: equal forward and reverse rates per volume when in equilibrium Boltzmann gives:  $\frac{\text{neutron } E_2 = m_{\vec{k}}^2}{2}$ 

$$\left(\frac{n}{p}\right)_{\text{equilib}} = \frac{g_n}{g_p} e^{-(E_2 - E_1)/T} = e^{-(m_n - m_n)/T} \left[\Delta E = \Delta m\right]$$

equilibrium requires *rate per nucleon*  $\Gamma \gg H$  *Hubble rate* lost at **freezeout**:

$$\Gamma \stackrel{\rm freeze}{=} H$$

 $n \leftrightarrow p$  freezeout:  $T_{\text{freeze}} \approx 1 \, \text{MeV}$ ,  $t_{\text{freeze}} \approx 1$  sec

# Interlude: Pair Annihilation

right after weak freezeout,  $T_{\gamma}$  drops below  $m_e = 0.511$  MeV pairs become nonrelativistic, annihilate:  $e^+e^- \rightarrow \gamma\gamma$ 

- mass energy  $\rightarrow$  back to radiation
- $\bullet$  small leftover amount of  $e^-$  to balance p charge
- ★ a sort of "heating" but really just restores relativistic energy  $T_{\gamma}$  never rises, but cooling is briefly slowed
- ★ since  $\nu$ s decoupled, don't receive pair energy cooler than photons thereafter can show (**PS3**):  $T_{\nu} = (4/11)^{1/3}T_{\gamma} = 0.714T_{\gamma}$ today, the (relativistic) cosmic neutrino backgrounds have  $T_{\nu,0} = 0.714T_{\gamma,0} = 1.95$  K

<sup>▶</sup> ★ if you can think of how to detect this *cosmic*  $\nu$  *background* let me know and we'll publish—you can even be second author!

### The Short but Interesting Life of a Neutron

(1) at 
$$T>T_f$$
,  $t\sim 1$  s  
 $n\leftrightarrow p$  rapid  
maintain  $n/p=(n/p)_{ ext{eq}}=e^{-\Delta m/T}$ 

(2) at  $T = T_f$ , freezeout fixes  $(n/p)_f = e^{-\Delta m/T_f} \simeq 1/6$ so *n* "mass fraction" is  $X_n = \frac{\rho_n}{\rho_B} = \frac{m_n n}{m_n n + m_p p} \approx \frac{n}{n+p} \approx 1/7$ (4)

(3) until nuclei form, free *n* decay:  $\dot{n} = -n/\tau_n$ , with  $\tau_n = 885.7 \pm 0.8$  s then mass fraction drops as

С

$$X_n = X_{n,i} e^{-\Delta t/\tau} \tag{5}$$

Q: why take this simple from?

# **Deuterium Bottleneck**

Build complex nuclei from n, pfirst step: deuterium production  $n + p \rightarrow d + \gamma$ www: BBN reaction network energy release  $Q = B(d) = E_{\gamma} = 2.22$  MeV: exothermic

reverse "photodissociation"  $d + \gamma \rightarrow n + p$  allowed but *endo*thermic

Naïvely: at  $T < T_f < Q$ , too cold to photo-dissociate But:  $n_{\gamma}/n_{\rm B} = 1/\eta \sim 10^9 \gg 1$  $\Rightarrow$  many photons per baryon  $\Rightarrow \langle E_{\gamma} \rangle < Q$ , but many photons have  $E_{\gamma} > Q$ D can't survive until  $T \ll Q!$ 

PS3: find when  $n_{\gamma}(>Q) \sim n_{\mathsf{B}}$ 

Q: How low to go?

σ

## **Nuclear Statistical Equilibrium**

For a non-relativistic species (Maxwell-Boltzmann):

$$n = \left(\frac{mkT}{2\pi\hbar^2}\right)^{3/2} e^{-(m-\mu)/T} \tag{6}$$

For  $n(p,\gamma)d$  in *chemical equilibrium*:  $\mu_n + \mu_p = \mu_d + \mu_\gamma = \mu_d$ , (since  $\mu_\gamma = 0$ ), so

$$\frac{n_n n_p}{n_d} = \left(\frac{(m_n m_p / m_d) kT}{2\pi \hbar^2}\right)^{3/2} e^{-(m_n + m_p - m_d)/T} \\
= \left(\frac{m_u kT/2}{2\pi \hbar^2}\right)^{3/2} e^{-B_D/T}$$
(7)

example of "nuclear statistical equilibrium" this example: Saha equation

write baryon fraction  $Y_i = n_i/n_B$  and  $n_B = \eta n_\gamma$ 

$$Y_d \sim Y_n Y_p \eta (T/m_u)^{3/2} e^{B_D/T}$$
(8)

Q: what is low-T behavior?

1

When  $Y_d \rightarrow 1$ : Nuke buildup starts

$$\ln Y_d \simeq B_D / T + \ln \eta + 3/2 \ln T / m_u \sim 0$$
 (9)

so (big hint for PS 3!)

$$T_D \simeq \frac{B_D}{\ln \eta^{-1}} \sim 0.07 \text{ MeV}$$
(10)

i.e., nuke rxns begin at  $T \simeq 10^9$  K Note:  $T_D \ll B_2$  since  $\eta \ll 1$ 

time  $t_d \sim 200 \text{ s} \rightarrow$  "the first 3 min"

between freezeout and  $T_D$ : free *n* decay: mass fraction  $X_n = X_{n,i}e^{-\Delta t/\tau} \simeq 0.12$ 

 $^{\infty}$  www: nuke network Q: where is flow direction? why?

Nuke reaction flow  $\rightarrow$  highest binding energy  $\rightarrow$  <sup>4</sup>He

almost all  $n \rightarrow {}^{4}$ He:  $n({}^{4}$ He)\_{after} = 1/2  $n(n)_{before}$   $Y_p = X({}^{4}$ He)  $\simeq 2(X_n)_{before} \simeq 0.24$   $\Rightarrow \sim 1/4$  of baryons into  ${}^{4}$ He,  $3/4 p \rightarrow$ H result weakly (log) dependent on  $\eta$ 

(11)

Robust prediction: large universal <sup>4</sup>He abundance

But nuke rxns also freeze out  $\Rightarrow n \rightarrow {}^{4}$ He conversion incomplete leftover traces of incomplete burning:

• D

Q

- <sup>3</sup>He (and <sup>3</sup>H $\rightarrow$ <sup>3</sup>He)
- <sup>7</sup>Li (and <sup>7</sup>Be $\rightarrow$ <sup>7</sup>Li)

trace abundances  $\leftrightarrow$  nuke freeze T

 $\Rightarrow$  strong  $n_B \propto \eta$  dependence

#### BBN theory: main result

- light element abundance predictions
- depend on baryon density  $\leftrightarrow \eta \leftrightarrow \Omega_{\text{baryon}}$

#### "Schramm Plot' '

10

Lite Elt Abundances vs  $\eta$ summarizes BBN theory predictions

www: Schramm plot

Note: no A > 7... Q: why not?

Why don't we go all the way to  ${}^{56}$ Fe? after all: most tightly bound  $\Rightarrow$  most favored thermodynamically (nuclear statistical equilibrium)

# Why no elements A > 7?

#### 1. Coulomb barrier

11

heavier products require heavier reactants which have higher charges

2. nuclear physics: "mass gaps" no stable nuclei have masses A = 5,8 $\rightarrow$  with just  $p \& {}^{4}$ He, can't overcome via 2-body rxs need 3-body rxns (e.g.,  $3\alpha \rightarrow {}^{12}$ C) to jump gaps but  $\rho$ , T too low

Stars *do* jump this gap, but only because have higher density a long time compared to BBN